

Basic crab cavities

ISA2003

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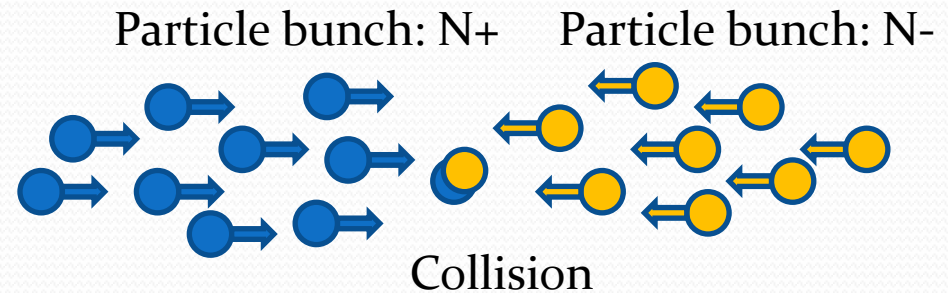
KEK

Basic crab cavities

- Luminosity of the collider accelerators and related topics
- RF cavities and parameters
- Pill box type crab cavities, mainly developed for KEKB
- Compact crab cavities, mainly developed for LHC

Colliding beams

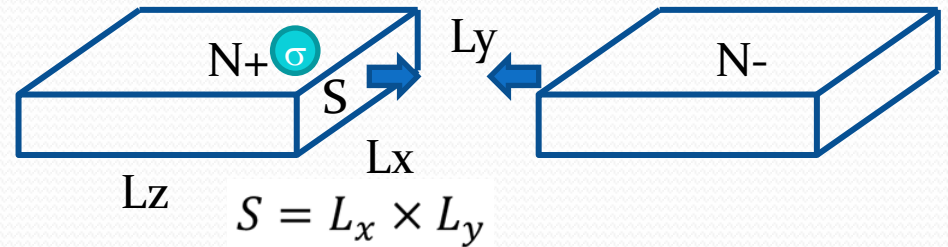
- We consider the luminosity of collider accelerators
- Luminosity is one of the most important parameters
- Number of physics interactions (events) per second = $L \times \sigma$
 - L : Luminosity
 - σ : Cross section
- L has to be increased to obtain high event rate



$$\frac{dN_{ev}}{dt} = L \cdot \sigma$$

Luminosity

- Calculate luminosity
 - Head on collision
 - Flat beams with uniform distribution
- # of events for a particle with a cross section of σ
- # of events for N_+ particles
- If those collisions occur f times per second
- Luminosity for flat beams



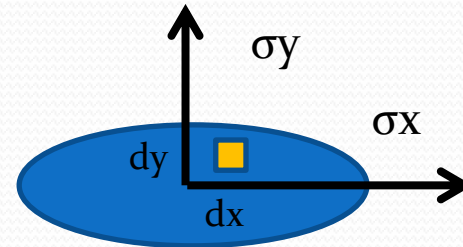
$$N_{ev} = N_- \sigma / S$$

$$N_{ev} = N_+ N_- \sigma / S$$

$$\frac{dN_{ev}}{dt} = f N_+ N_- \sigma / S$$

$$L = \frac{N_+ N_-}{L_x L_y} f$$

Luminosity



- In case of Gaussian distribution we have to modify the previous equation

- Same σ_x , σ_y for both beams

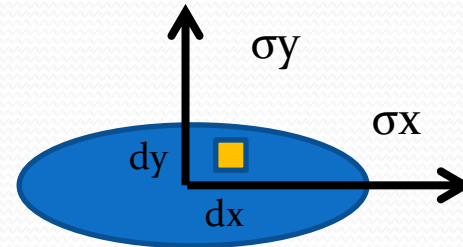
$$L = \frac{N_+ N_-}{4\pi\sigma_x\sigma_y} f$$

$$L = \frac{N_+ N_-}{L_x L_y} f \quad \begin{array}{l} L_x \rightarrow dx \\ L_y \rightarrow dy \end{array}$$

$$N_+ \rightarrow N_+ \frac{1}{\sqrt{2\pi}\sigma_{x+}} e^{-\frac{x^2}{2\sigma_{x+}^2}} \frac{1}{\sqrt{2\pi}\sigma_{y+}} e^{-\frac{y^2}{2\sigma_{y+}^2}} dx dy$$

$$N_- \rightarrow N_- \frac{1}{\sqrt{2\pi}\sigma_{x-}} e^{-\frac{x^2}{2\sigma_{x-}^2}} \frac{1}{\sqrt{2\pi}\sigma_{y-}} e^{-\frac{y^2}{2\sigma_{y-}^2}} dx dy$$

Luminosity



- In case of Gaussian distribution
 - Same σ_x , σ_y
 - But the other beam has offset Δx
- Reduction factor

$$L = \frac{N_+ N_-}{4\pi\sigma_x\sigma_y} f$$

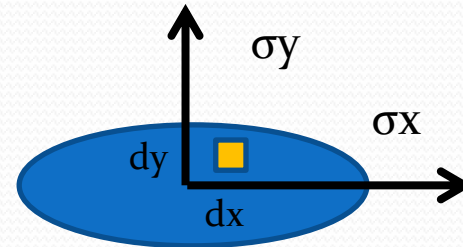
$$N_- \rightarrow N_- \frac{1}{\sqrt{2\pi}\sigma_{x-}} e^{-\frac{(x-\Delta x)^2}{2\sigma_{x-}^2}} \frac{1}{\sqrt{2\pi}\sigma_{y-}} e^{-\frac{y^2}{2\sigma_{y-}^2}} dx dy$$

$$R_{\Delta x} = e^{-\frac{\Delta x^2}{4\sigma_x^2}}$$

Luminosity

- In case of Gaussian distribution
 - Different σ_x , σ_y
 - Define Σ_x and Σ_y

$$L = \frac{N_+ N_-}{2\pi \Sigma_x \Sigma_y} f$$



$$\Sigma_x = \sqrt{\sigma_{x+}^2 + \sigma_{x-}^2}$$

$$\Sigma_y = \sqrt{\sigma_{y+}^2 + \sigma_{y-}^2}$$

Finite angle crossing

- To increase luminosity
 - Increase beam bunches
 - When bunches increased parasitic collisions occur
- To avoid parasitic collisions
 - Beam crossing with finite angle: θ_c
 - Make beam separation simple
 - Reduction of luminosity
 - Geometrical reduction factor R_L
 - May introduce coupling between synchrotron and betatron oscillations in circular colliders



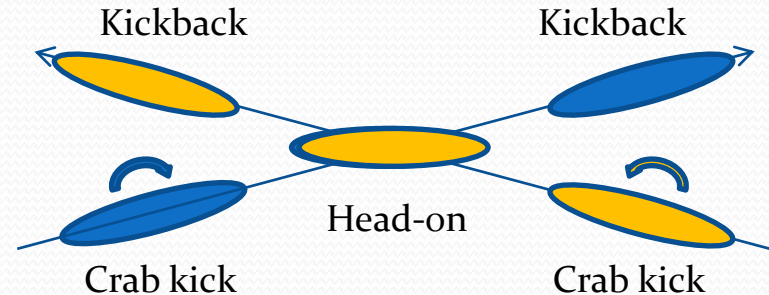
$$R_L = \frac{1}{\sqrt{1 + \phi_P^2}}$$

$$\phi_P = \frac{\sigma_z}{\sigma_x} \tan \theta_c / 2$$

Piwinski angle

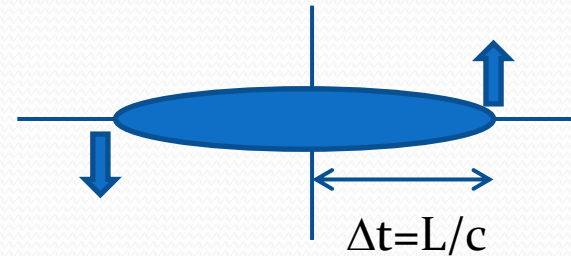
Crab crossing

- Bunch rotation at the IP
 - Horizontal kick at the bunch head
 - Horizontal kick at the bunch tail but opposite direction
 - No kick at the bunch center
- Bunch rotates so that it collides head-on at IP
 - Crab crossing
- Crab cavity kick the bunch with a phase advance of $\pi/2$
- Another crab cavity kicks back the bunch after IP



Required kick voltage

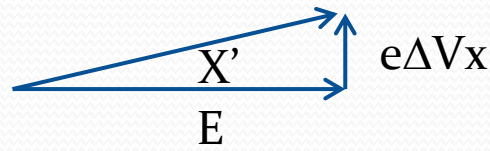
- Kick voltage given to the bunch head
 - L: length from the bunch center
 - c: speed of light
 - ΔV : kick voltage
 - No kick to the bunch center
 - Same kick voltage to the bunch tail, but opposite direction



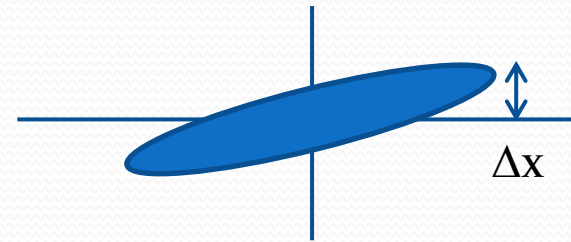
$$\Delta V_x = V_x \sin \omega_{RF} \Delta t \sim V_x \omega_{RF} \frac{L}{c}$$

Required kick voltage

- Kick angle, x'
- Offset at the IP
 - m_{12} : transfer matrix component of the betatron oscillation
 - $\Delta\phi$: phase advance ($\sin\Delta\phi=1$)
- Bunch rotation angle at the IP
 - Half crossing angle
- Required kick voltage



$$x' = \frac{e\Delta V_x}{E} = \frac{\Delta V_x}{E/e}$$



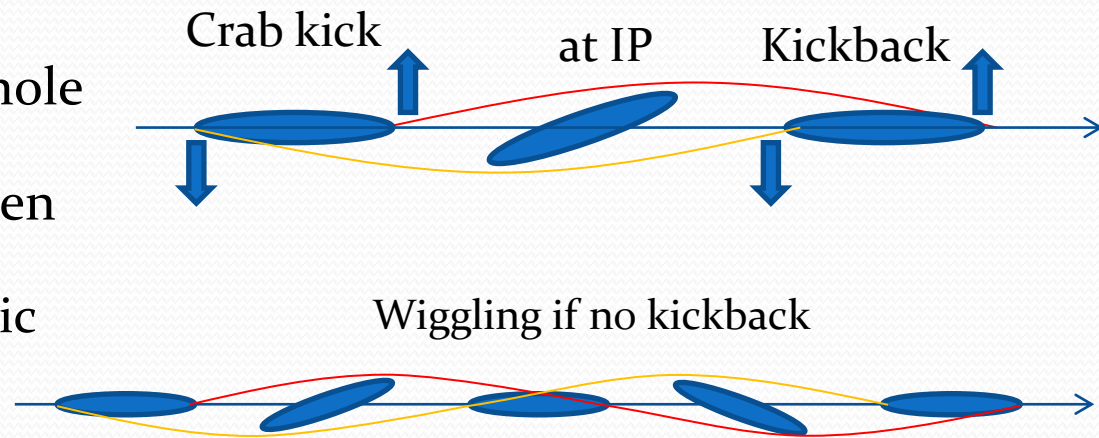
$$\Delta x = m_{12}x' = \sqrt{\beta_{crab}\beta_{IP}} \sin \Delta\phi x'$$

$$\tan \theta = \frac{\Delta x}{L} = \sqrt{\beta_{crab}\beta_{IP}} \frac{V_x \omega_{RF}}{cE/e}$$

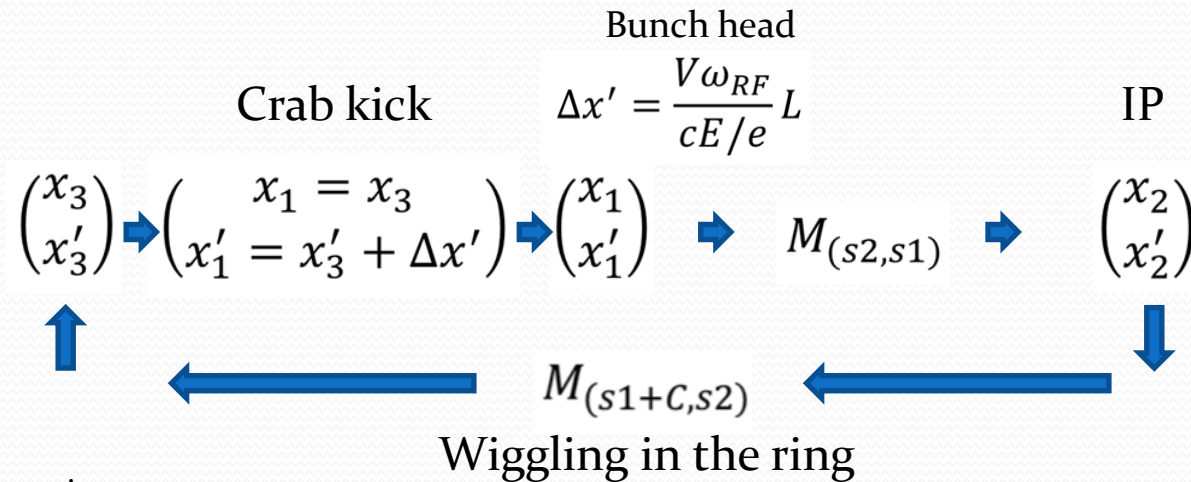
$$V_x = \frac{cE/e}{\sqrt{\beta_{crab}\beta_{IP}} \omega_{RF}} \theta$$

Single crab cavity scheme

- One crab cavity in the ring
- The beam wiggles in the whole ring
- Cavity location can be chosen anywhere in the ring
 - near the existing cryogenic facility
- KEKB applied this scheme
 - Two crab cavities fabricated and installed in the HER and LER ring
 - The location is 1 km away from the IP, where a cryogenic facility was operating for the SRF accelerating cavities
 - Great cost reduction



Required kick voltage in single crab scheme



Coordinate transformation

$$\begin{pmatrix} x_2 \\ x'_2 \end{pmatrix} = M_{(s2,s1)} \begin{pmatrix} x_1 \\ x'_1 \end{pmatrix}$$

$$\begin{pmatrix} x_3 \\ x'_3 \end{pmatrix} = M_{(s1+C,s2)} \begin{pmatrix} x_2 \\ x'_2 \end{pmatrix}$$

Transfer matrix

$$M_{(s2,s1)} = \begin{pmatrix} \sqrt{\frac{\beta_2}{\beta_1}} (\cos \psi_{21} + \alpha_1 \sin \psi_{21}) & \sqrt{\beta_2 \beta_1} \sin \psi_{21} \\ \frac{1}{\sqrt{\beta_2 \beta_1}} (-(\alpha_2 - \alpha_1) \cos \psi_{21} - (1 + \alpha_2 \alpha_1) \sin \psi_{21}) & \sqrt{\frac{\beta_2}{\beta_1}} (\cos \psi_{21} - \alpha_1 \sin \psi_{21}) \end{pmatrix}$$

Bunch rotation angle at IP

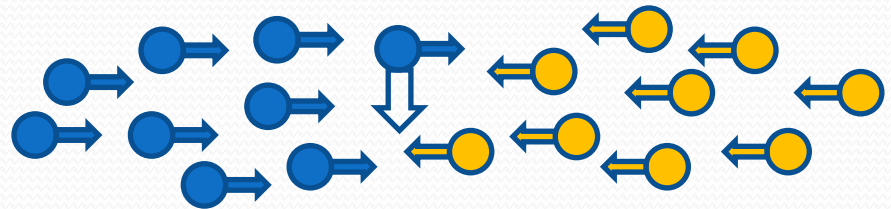
$$\tan \frac{\theta_c}{2} = \sqrt{\beta_2 \beta_1} \frac{\cos \left(\frac{\psi_c}{2} - \psi_{21} \right)}{2 \sin \frac{\psi_c}{2}} \frac{V\omega_{RF}}{cE/e}$$

Phase advance in the ring: $\psi_c \sim \pi$

Phase advance at the crab: $\psi_{21} = \pi/2$

Beam-beam parameter

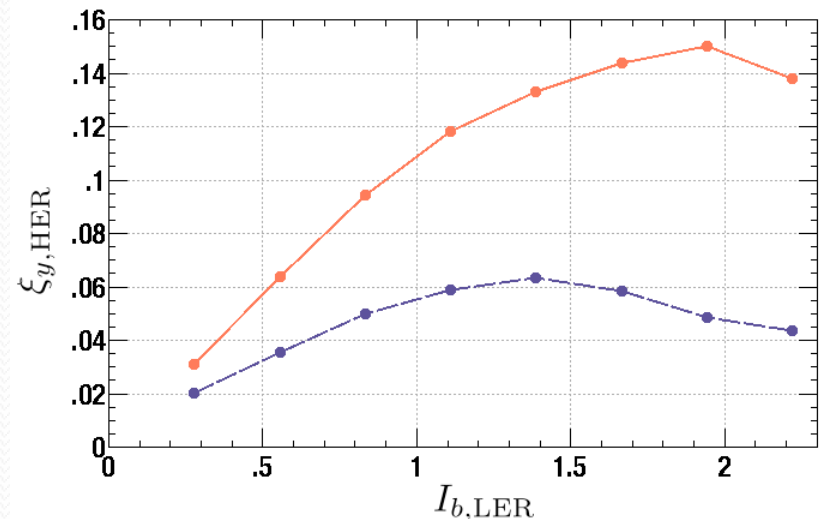
- Beam-beam kick
 - Electromagnetic force of electrons kicks positrons
 - Kicking force is approximately proportional to the distance from the center
- Tune shift
 - This force is like a quadrupole magnet
 - Tune shifts due to this kick
 - $\Delta\mu$: ξ parameter
 - Beam-beam parameter
- Beam-beam limit
 - If we increase N , σ increases
 - Beam-beam parameter has a limit
 - Beam-beam limit
 - Luminosity has a limit
- Beam-beam simulation predicted ξ could be increased more than twice in the crab crossing



$$\xi_{x,y\pm} = \frac{r_e N_{\mp}}{2\pi\gamma_{\pm}} \frac{\beta_{x,y\pm}}{\sigma_{x,y\mp}\sigma_{y\mp}}$$

$$L = \frac{N_{\pm}}{2r_e} f \gamma_{\pm} \frac{\xi_{y\pm}}{\beta_{y\pm}}$$

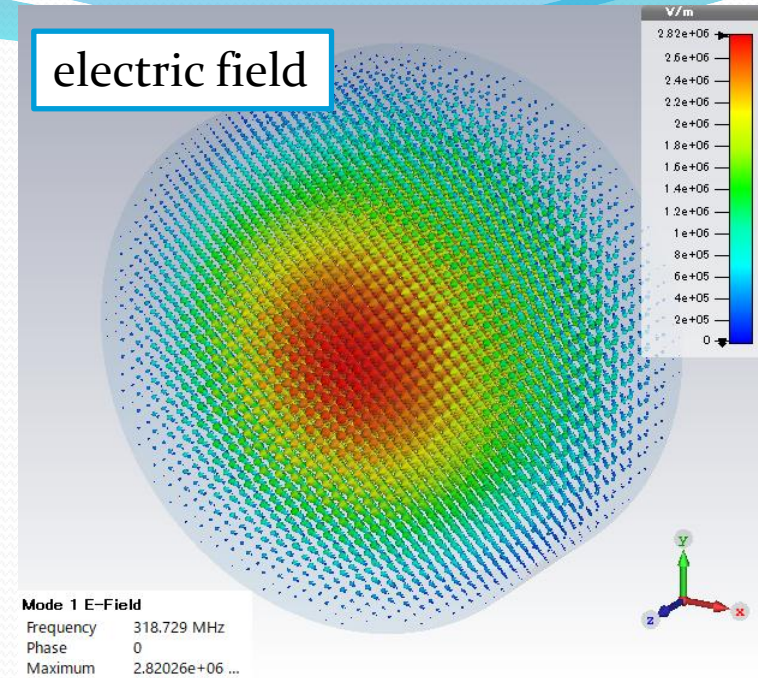
r_e : Classical electron radius
 γ : Lorentz factor



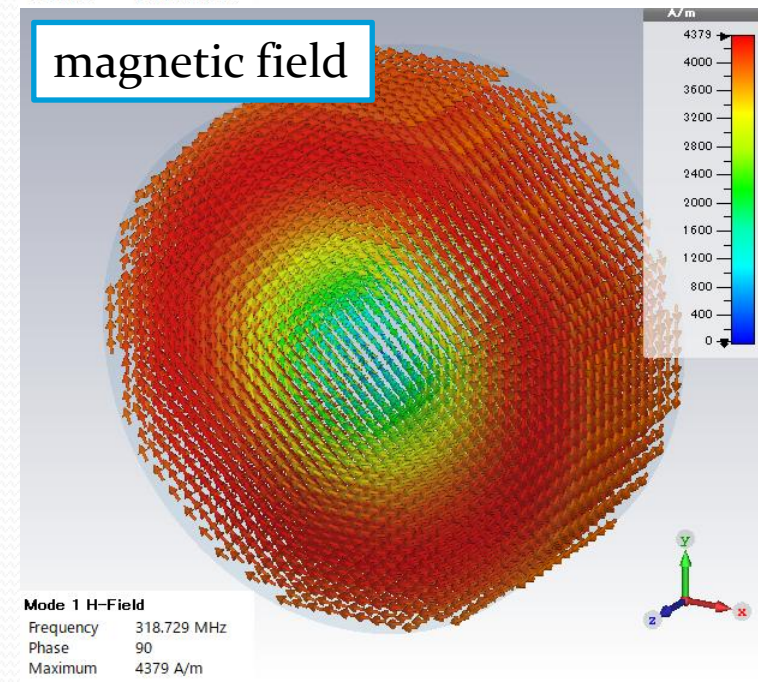
RF cavity

- Particle acceleration longitudinally or transversely
- Time varying electric and magnetic fields in a cavity
 - Resonate with the resonant frequency
 - Create large electric and magnetic fields
 - Those fields accelerate charged particles

electric field



magnetic field



RF cavity

- Resonant frequency and electro magnetic fields can be obtained from;
 - Maxell's equations
 - Boundary conditions
- Electro-magnetic fields decay in time due to the heating loss on the boundary
 - The ratio of the loss in one cycle P_c to stored energy U characterizes cavity performance
 - Quality factor: $Q = \omega U / P_c$

$$\text{Div} B = 0$$

$$\text{Div} E = 0$$

$$\text{Rot} E = - \frac{\partial B}{\partial t}$$

$$\text{Rot} H = \frac{\partial D}{\partial t}$$

$$Q = \frac{\omega U}{P_c}$$

Beam acceleration

- Lorentz force
 - Charged particle with speed of light enters into a cavity
 - Charged particles obtain acceleration: Lorentz force
- We obtain the accelerating voltage when the cavity has longitudinal electric field
- Accelerating voltage V_c
 - The EM fields vary in time
 - We have to take the travelling time of a particle into account
 - Transit time factor: T
 - The ratio of V_c to the maximum voltage

$$F = e(E + v \times B)$$

$$F_z = eE_z$$

$$V_c = \int_0^d E_z e^{i\omega t} dz = \int_0^d E_z e^{ikz} dz$$

$$V_c = \left| \int_0^d E_z e^{ikz} dz \right|$$

$$T = \frac{\left| \int_0^d E_z e^{ikz} dz \right|}{\left| \int_0^d E_z dz \right|}$$

Shunt impedance

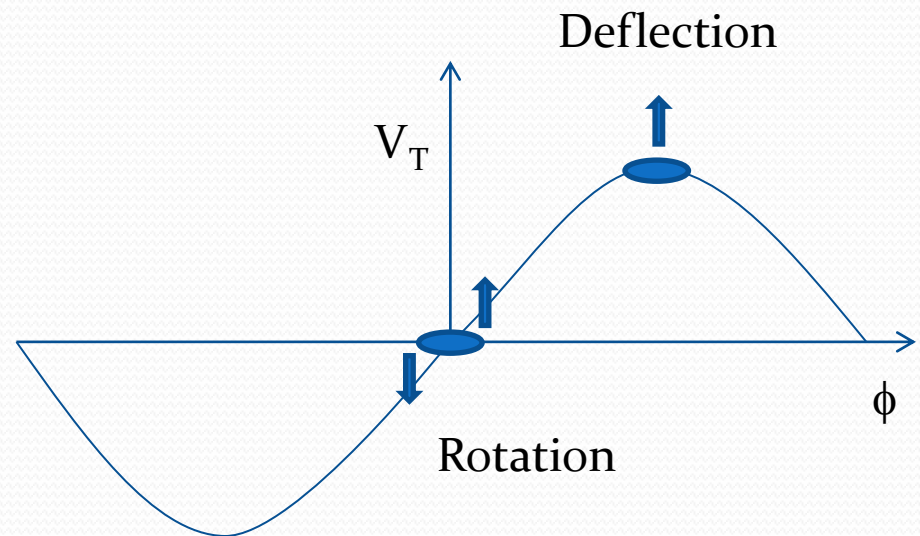
- Shunt impedance
 - Important parameter
 - Ratio of V_c^2 to P_c
 - Efficiency of acceleration
 - In the lumped circuit theory; $R = V_c^2 / P_c / 2$
- R/Q
 - Another important parameter

$$R = \frac{V_c^2}{P_c}$$

$$R/Q = \frac{V_c^2}{\omega U}$$

Direction of acceleration

- Longitudinal acceleration
 - Beam acceleration
- Transverse acceleration
 - Beam deflection
 - Whole bunch deflects
 - Bunch rotation
 - Bunch head deflects
 - Bunch tail deflects, but in the opposite direction
 - Bunch center does not deflect
 - Crabbing



How to kick the beam

- Lorentz force
 - Particle with $v=c$ along z-axis
- The force accelerates particles
 - Z-comp: acceleration
 - X-comp: horizontal kick
 - Y-comp: vertical kick
- Horizontal kick voltage

$$F = e(E + v \times B)$$

$$v = ce_z = (0, 0, c)$$

$$F_x = e(E_x - cB_y)$$

$$F_y = e(E_y + cB_x)$$

$$F_z = eE_z$$

$$V_x = \int_0^d (E_x - cB_y) e^{ikz} dz$$

Panofsky-Wenzel theorem

- Maxwell's eq: $\text{Rot}E = -\frac{\partial B}{\partial t} = -i\omega B \rightarrow B_y = \pi \frac{i}{\omega} \left(\frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} \right)$

- V_x can be given by E only

$$V_x = \int_0^d \left(E_x - i \frac{c}{\omega} \frac{\partial E_x}{\partial z} + i \frac{c}{\omega} \frac{\partial E_z}{\partial x} \right) e^{ikz} dz$$

- Note that the first 2 terms are combined

$$\left(E_x - i \frac{c}{\omega} \frac{\partial E_x}{\partial z} \right) e^{ikz} = -i \frac{c}{\omega} \frac{d}{dz} (E_x e^{ikz})$$

- V_x is given by the gradient of V_z

$$V_x = \int_0^d i \frac{c}{\omega} \frac{\partial E_z}{\partial x} e^{ikz} dz = \frac{i}{k} \frac{\partial V_z}{\partial x}$$

- The TE mode does not kick the particle because $E_z=0$

Kick voltage

- Kick voltage
 - Horizontal kick
 - Vertical kick
- Shunt impedance
 - T stands for x or y
- R/Q

$$V_x = \frac{i}{k} \frac{\partial V_z}{\partial x}$$

$$V_y = \frac{i}{k} \frac{\partial V_z}{\partial y}$$

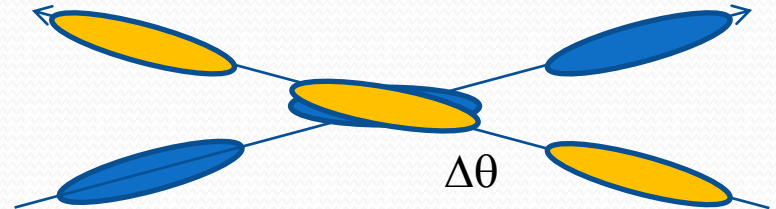
$$R_T = \frac{V_T^2}{P_c}$$

$$R/Q = \frac{V_T^2}{\omega U}$$

Voltage stability

- Voltage error: ΔV
- Bunch rotation angle error: $\Delta\theta$
 - Proportional to V
- Geometrical reduction factor

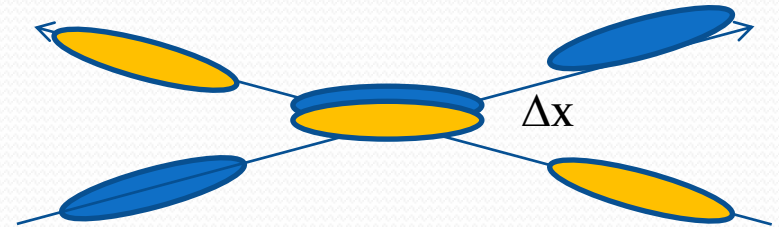
$$R_L = \frac{1}{\sqrt{1 + \phi_P^2}}$$



$$\Delta\theta = \sqrt{\beta_{crab}\beta_{IP}} \frac{\omega_{RF}}{cE/e} \Delta V$$

Phase stability

- Phase error of one crab cavity: $\Delta\phi$
- Transverse kick of the bunch center
 - Proportional to $\Delta\phi$
- Bunch position offset at IP: Δx
- Geometrical reduction factor
- Absolute phase error of two crab cavities has no geometrical reduction
 - Both beams have the same offset at IP
 - Because $\Delta x=0$



$$\Delta x = \sqrt{\beta_{crab}\beta_{IP}} \frac{V}{E/e} \Delta\phi$$

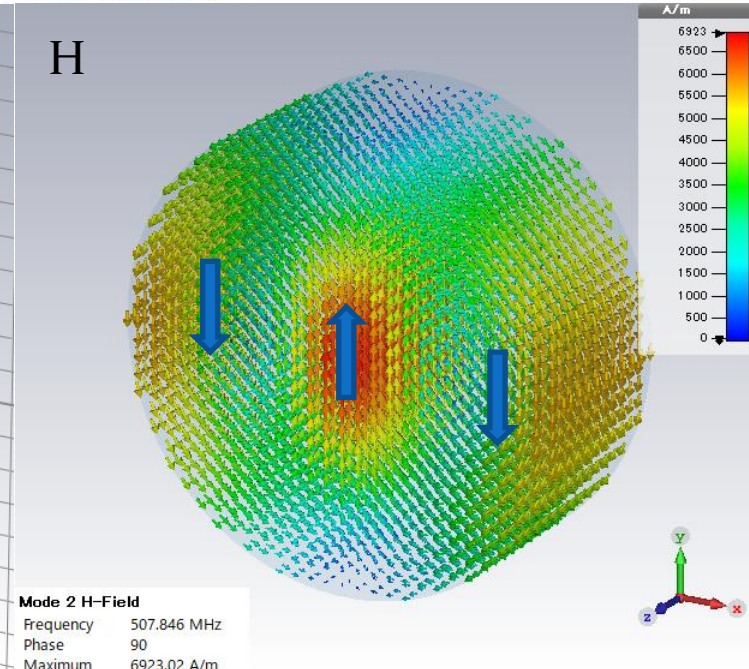
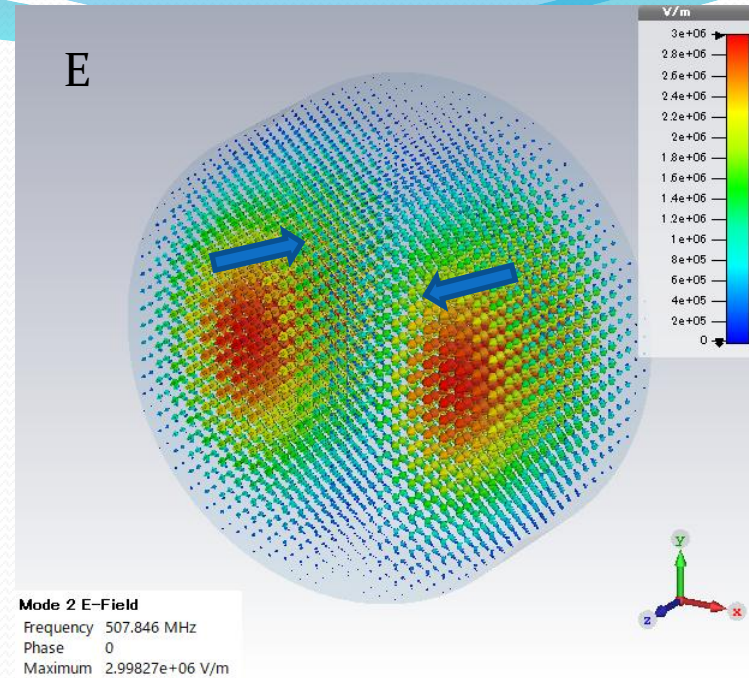
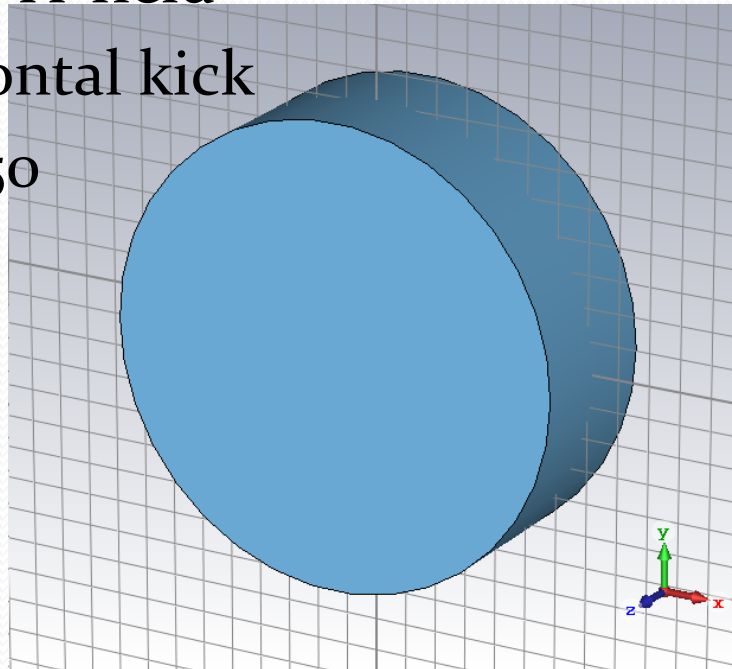
$$R_{\Delta x} = e^{-\frac{\Delta x^2}{4\sigma_x^2}}$$

Cavity type

- Pill box
- Rectangular cavity
- Coaxial cavity

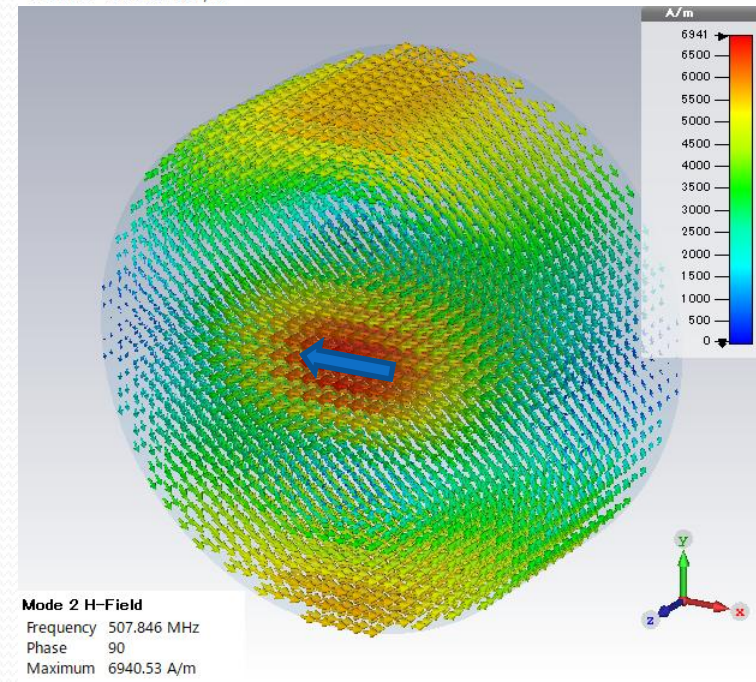
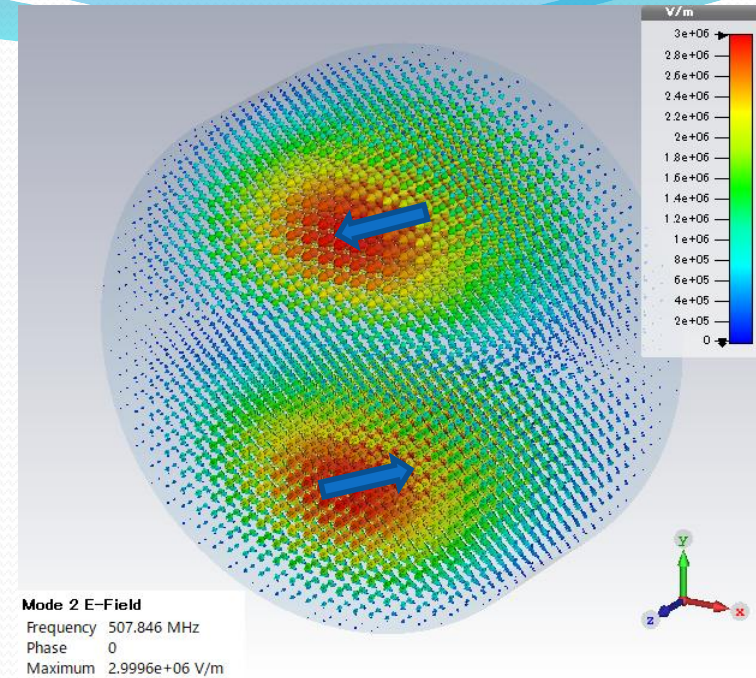
Crab cavity shape

- Pill box type
- TM₁₁₀ dipole mode
 - 500MHz
 - E- and H-field
 - Horizontal kick
 - R/Q~50



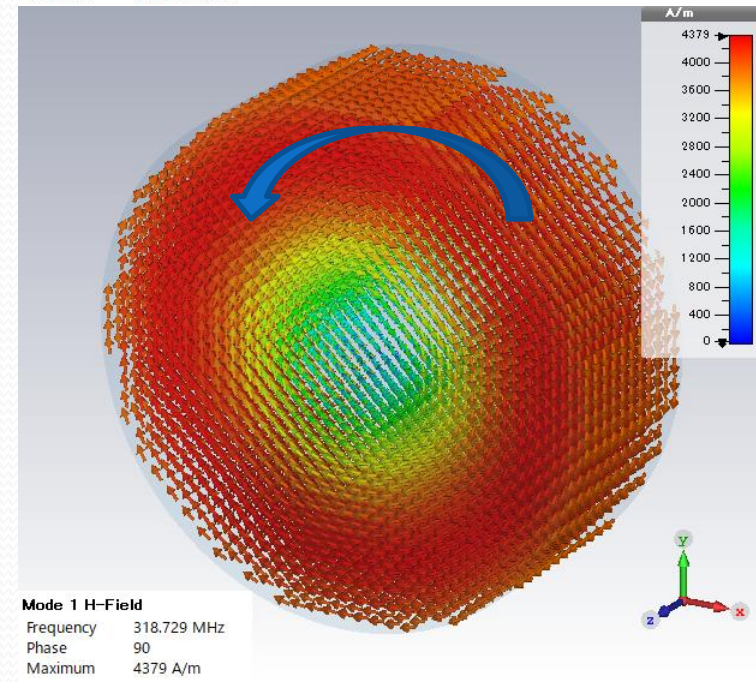
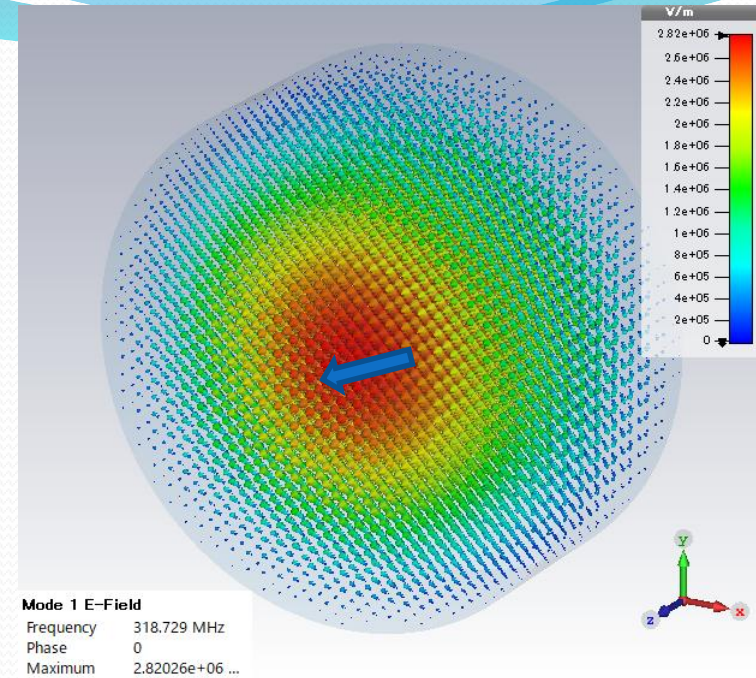
SOM

- Same order mode
 - Unwanted polarization
 - Vertical kick
 - Same frequency



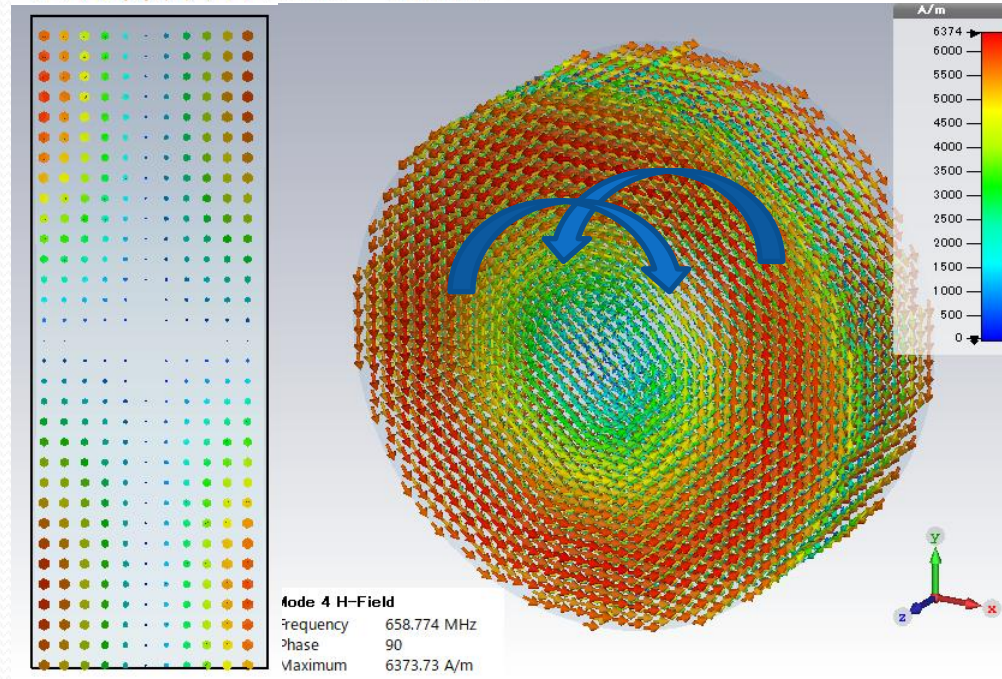
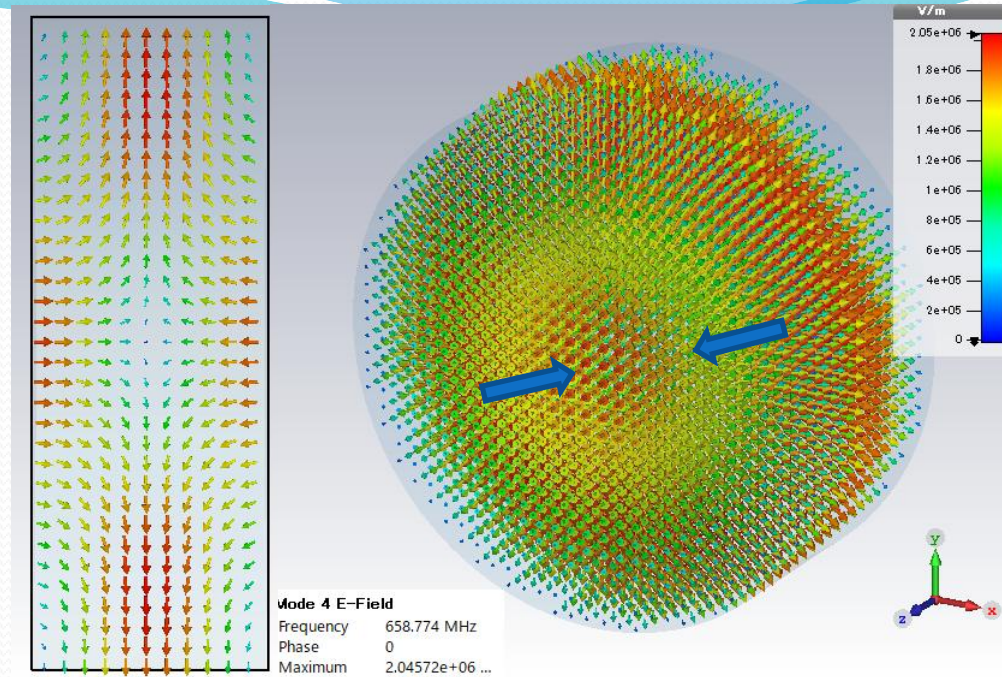
LOM

- Lowest order mode
 - Fundamental mode
- TM₀₁₀ monopole mode
 - Lower frequency: 320 MHz
 - Accelerating mode
 - R/Q~100
 - Need to be heavily damped



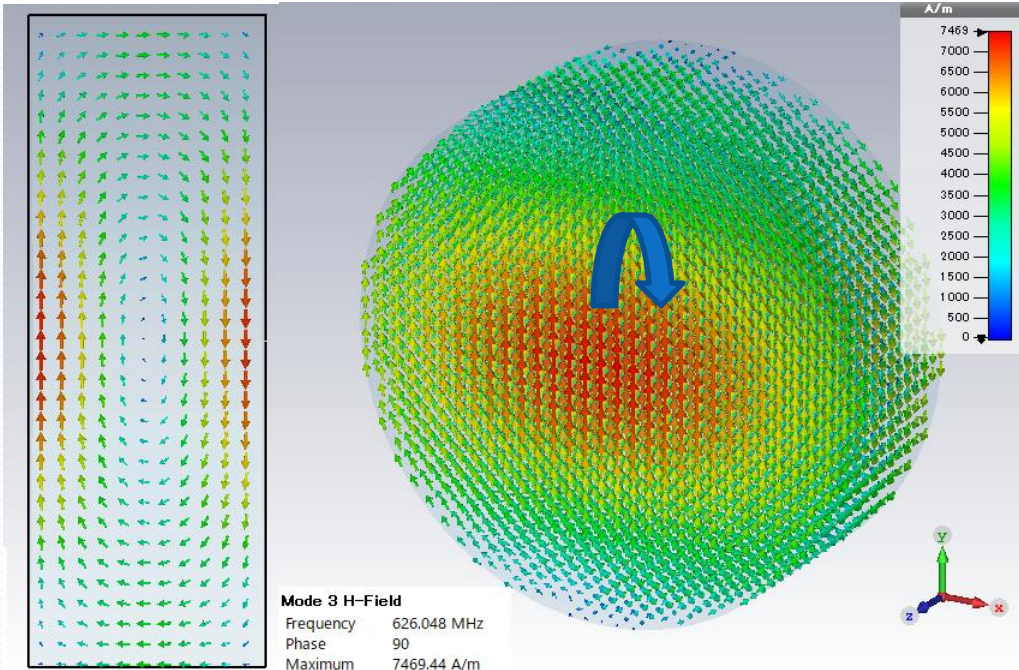
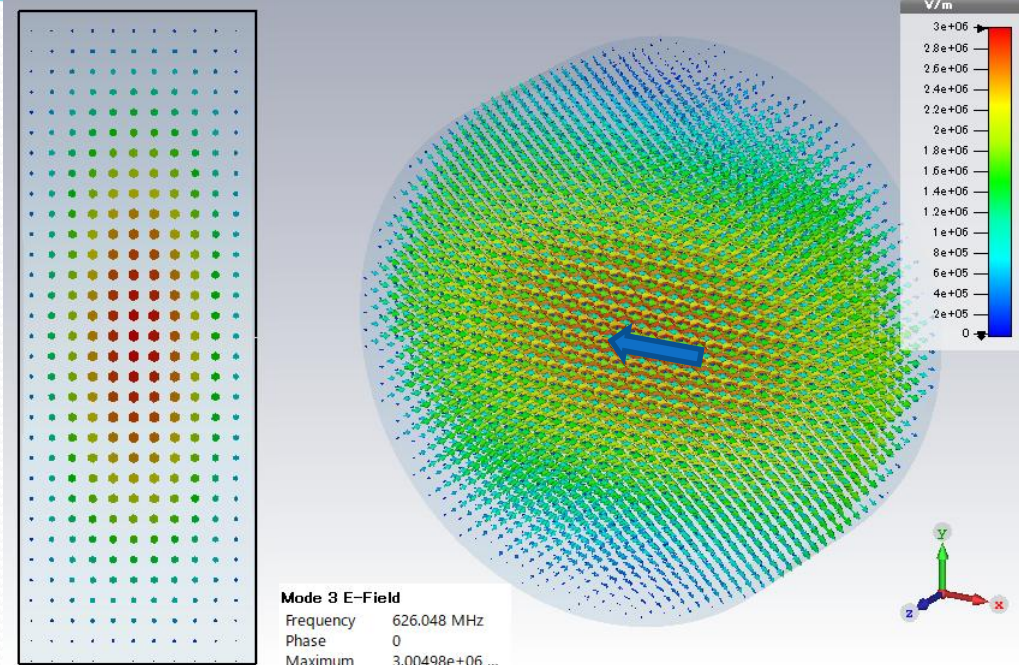
HOM

- Higher order mode
- TM₀₁₁ monopole mode
- Accelerating mode
- Should be heavily dumped



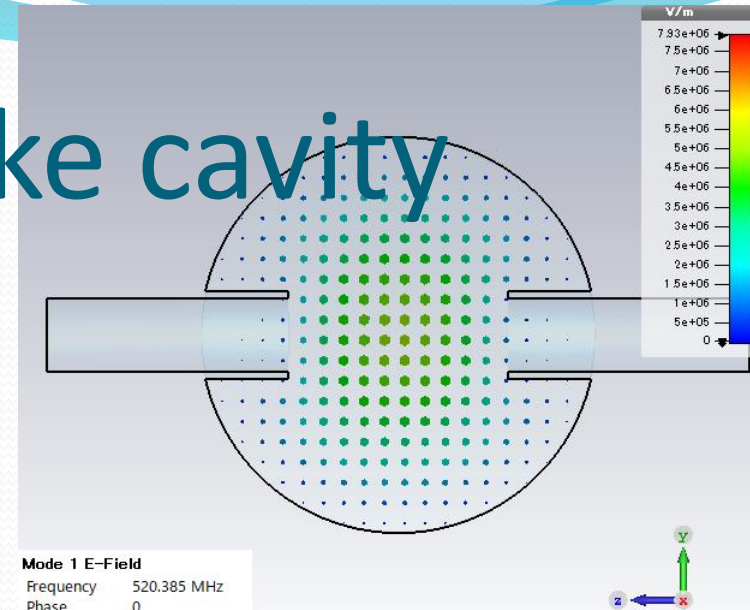
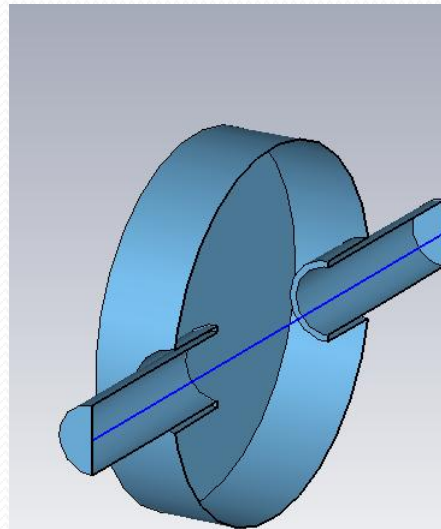
HOM, Crab cavity?

- TE₁₁₁ dipole mode
- Crab cavity?
- Answer is No!
- Magnetic kick cancels electric kick
 - No kick when $E_z=0$

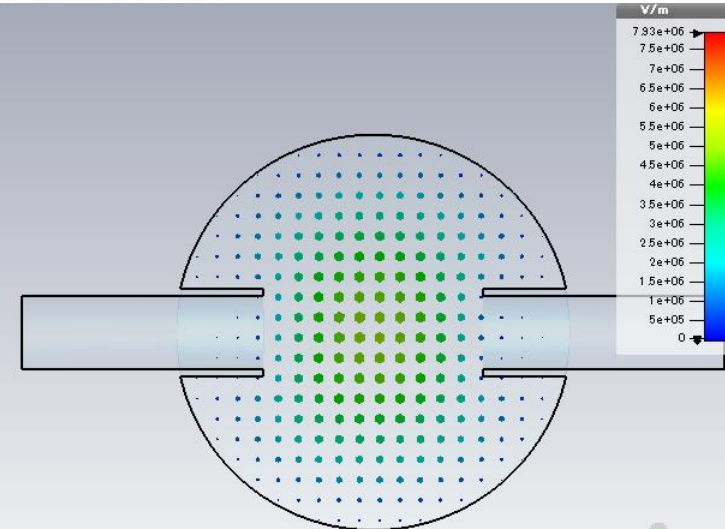


We can design TE-like cavity

- TE cavity with beam pipe
 - Frequency: 500 MHz
- If the beam pipes shield the magnetic field
- We can obtain a horizontal kick
 - $R/Q \sim 40$



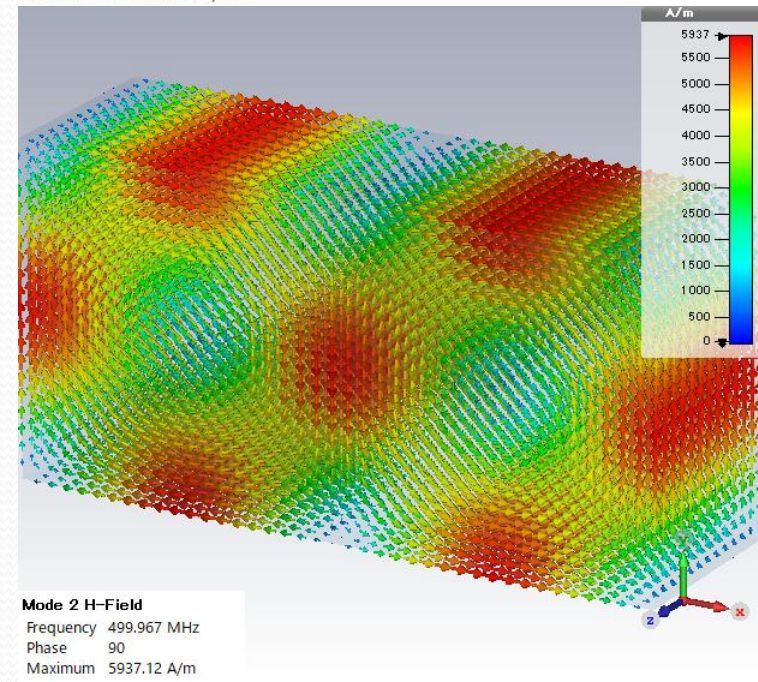
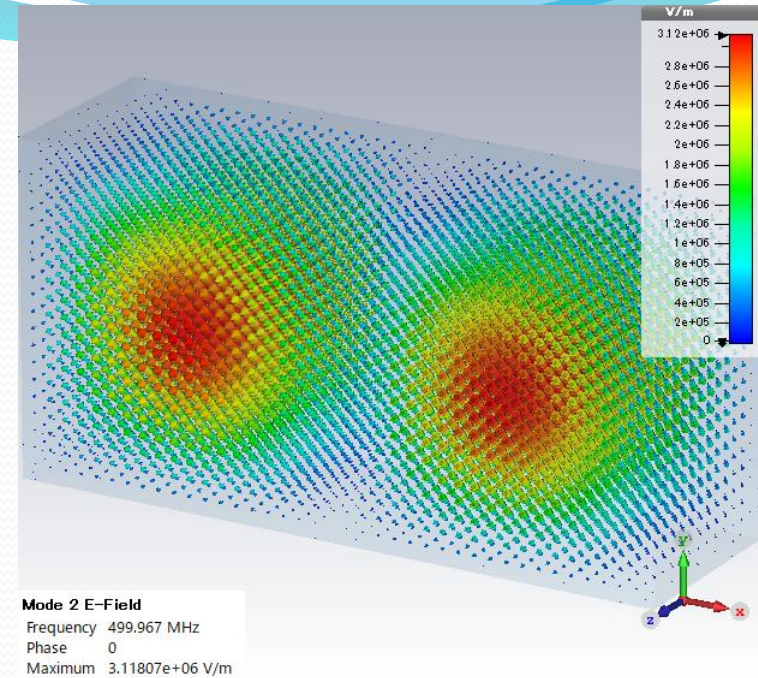
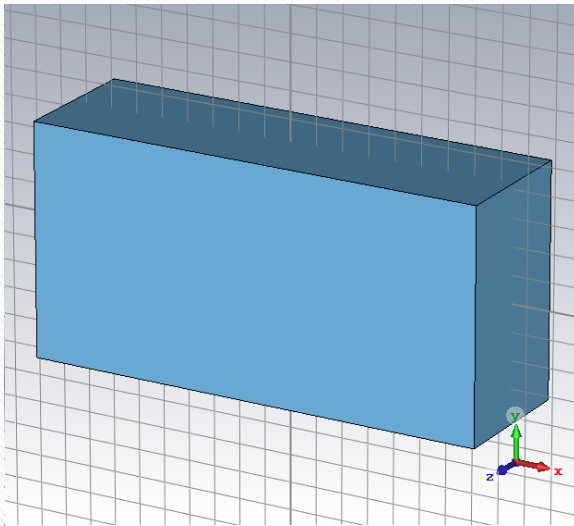
Mode 1 E-Field
Frequency 520.385 MHz
Phase 0
Cross section A
Cutplane at X 0.000
Maximum 4.96199e+06 ...



Mode 1 E-Field
Frequency 520.385 MHz
Phase 0
Cross section A
Cutplane at X 0.000
Maximum 4.96199e+06 ...

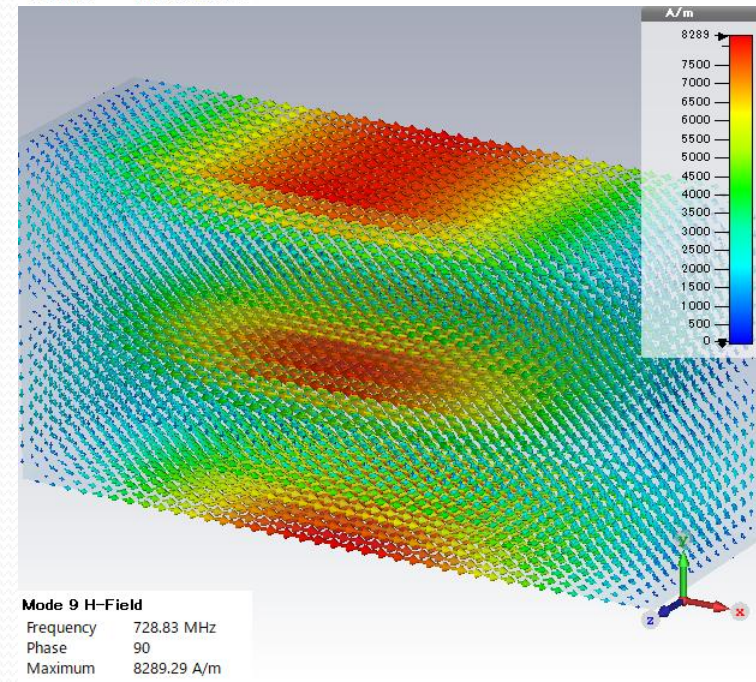
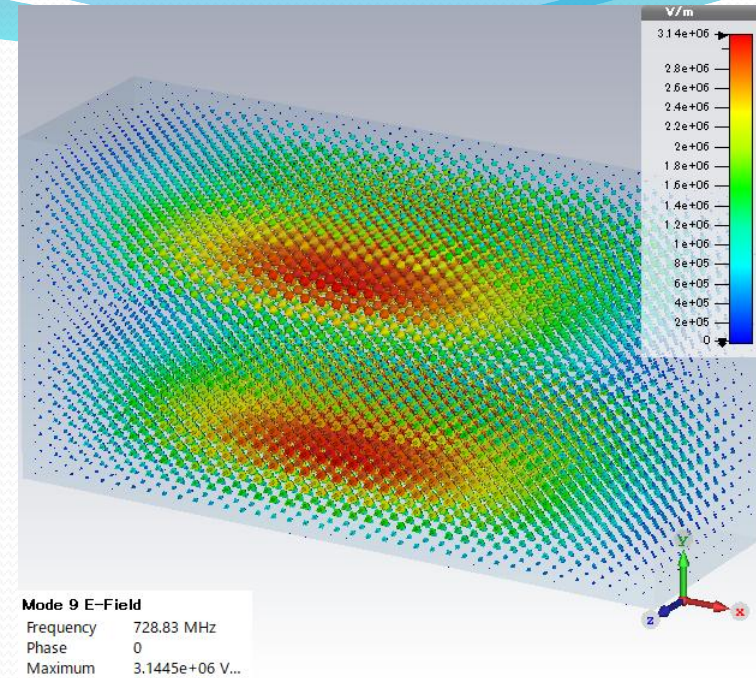
Rectangular cell

- Crabbing mode, TM₂₁₀
- 500MHz



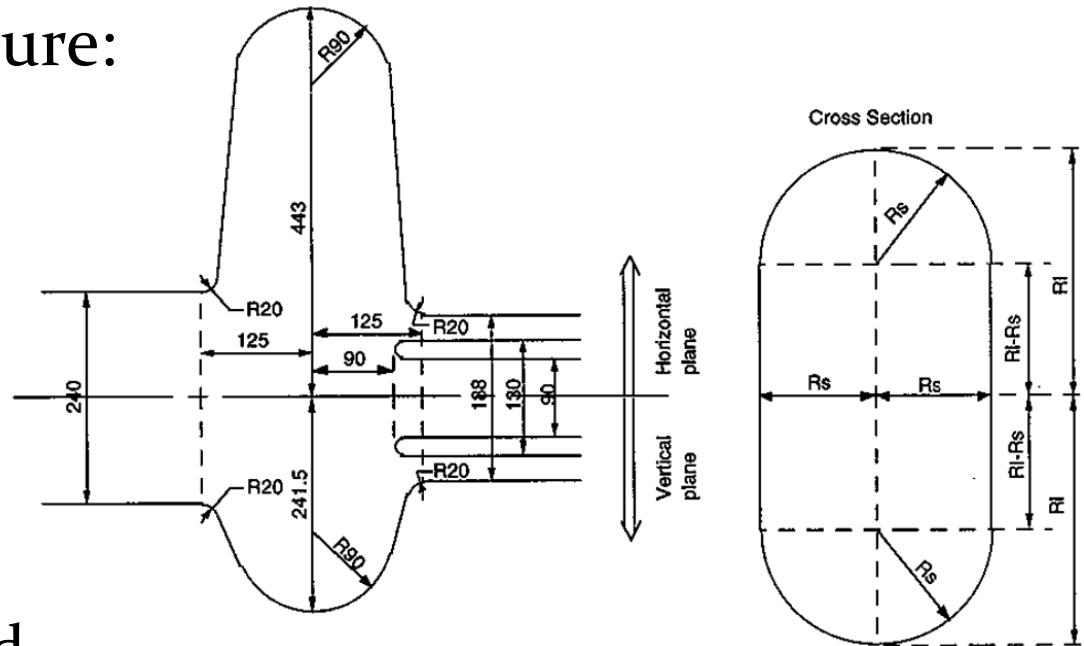
SOM

- TM₁₂₀
- Vertical kicking mode
- Higher frequency than the crabbing mode
- We can separate vertical kicking mode from horizontal kick mode



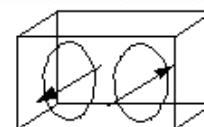
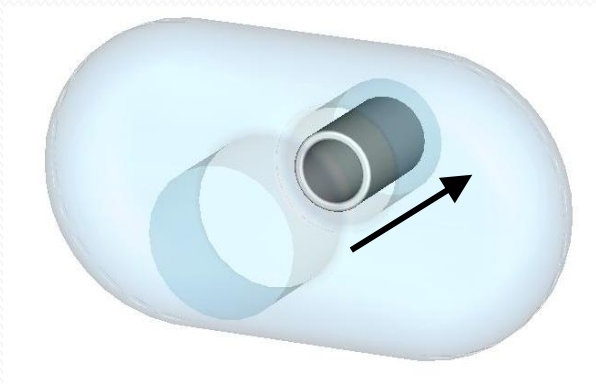
Crab cavity cell for KEKB

- Required voltage: 1.4MV
- Operating temperature: 4.4K
- Squashed cell
- Crabbing mode
 - TM₂₁₀-like mode
 - 509 MHz
- Coaxial coupler
 - For LOM, SOM and HOMs damping



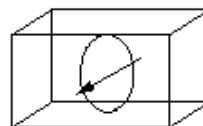
LOM, SOM and HOMs

- Coaxial coupler propagates these modes
- TEM: No cut-off (monopole)
- TE₁₁: $f_c = 600$ MHz (dipole)

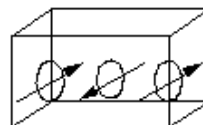


TM₂₁₀
Crab Mode

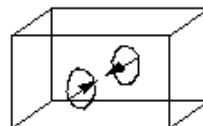
Monopole



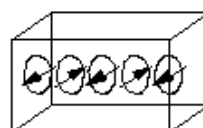
TM₁₁₀
LOM



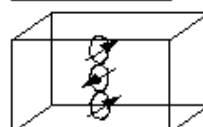
TM₃₁₀



TM₁₁₁

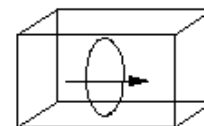


TM₅₁₀

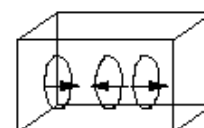


TM₁₃₀

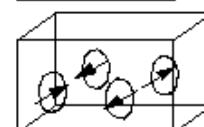
Dipole
Horizontal Polarization



TE₀₁₁

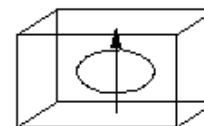


TE₂₁₁

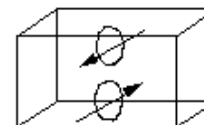


TM₂₁₁

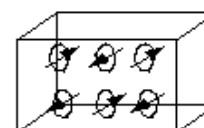
Dipole
Vertical Polarization



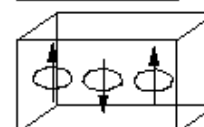
TE₁₀₁



TM₁₂₀



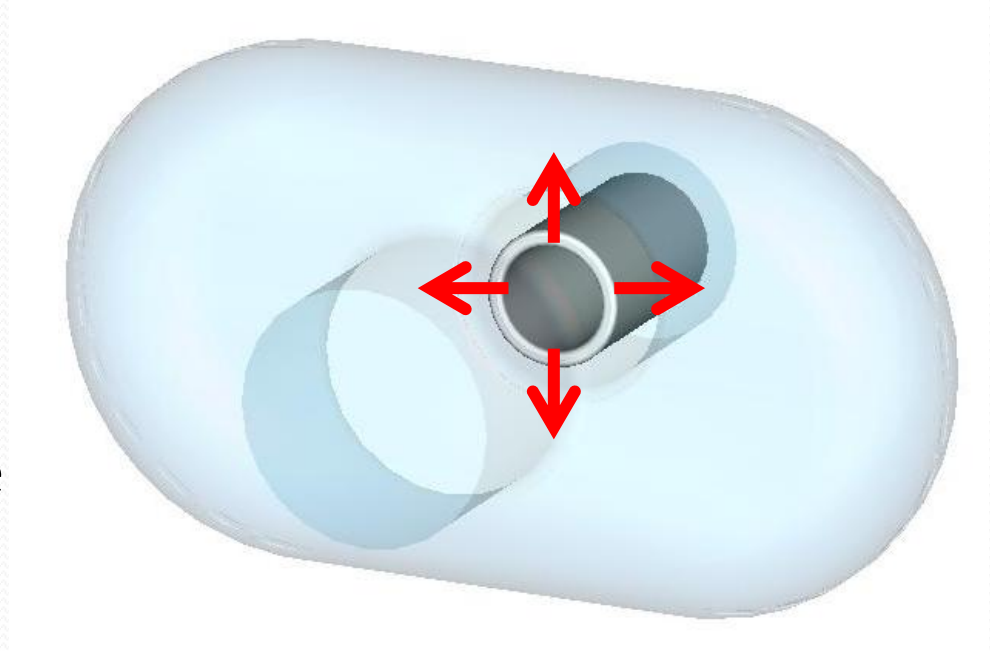
TE₃₂₀



TM₃₀₁

LOM damping

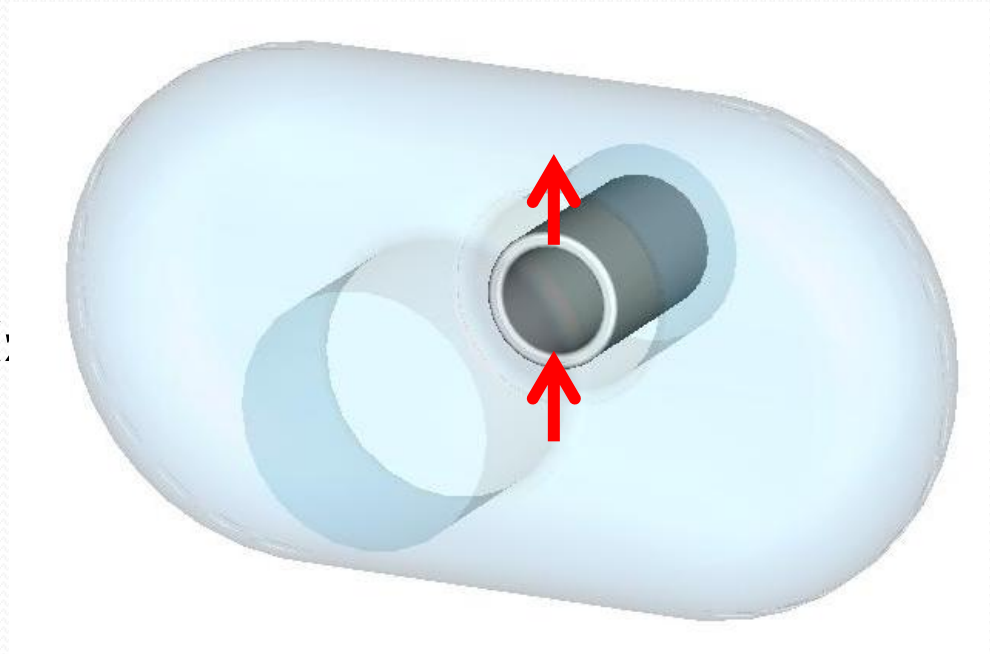
- Lowest order mode
 - Fundamental mode
 - Accelerating mode
 - $R/Q \sim 100$
- LOM can propagate as the TEM mode in the coaxial coupler
- Coaxial ferrite RF absorber at the end of this coupler heavily damps this mode
 - $Q \sim 100$



TEM mode in the coaxial coupler

SOM damping

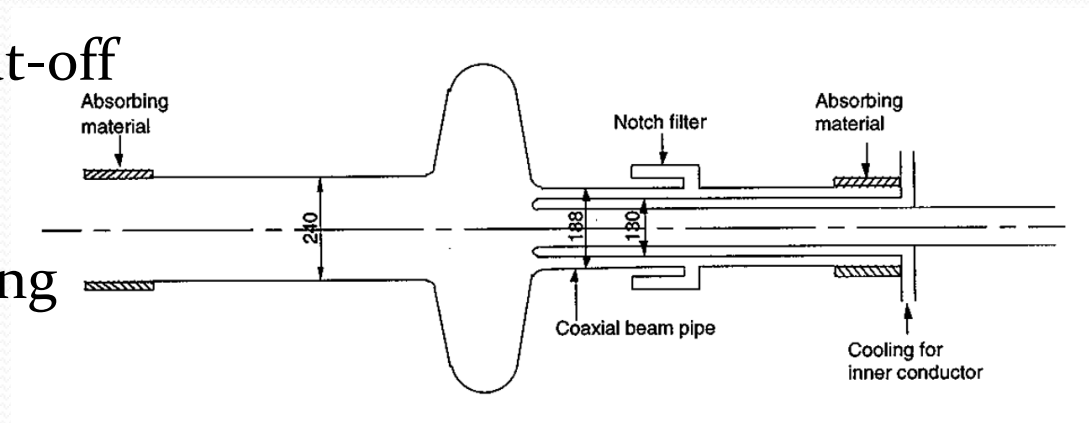
- Same order mode
 - Another polarization of crabbing mode
 - Frequency above 600 MHz by squeezing cavity cell
- SOM can propagate as the dipole mode (TE_{11}) in the coaxial coupler
 - The cut-off frequency is 600 MHz
 - Above the crabbing mode
 - Blow the SOM mode



TE mode in the coaxial coupler

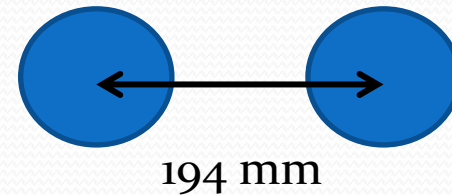
HOM damping

- Coaxial coupler and coaxial ferrite absorber
 - Monopole modes
 - Dipole mode above cut-off
 - Heavily damp those modes
 - Notch filter for crabbing mode rejection
- Large beam pipe and ferrite RF absorbers
 - Monopole > 900 MHz
 - Dipole > 750 MHz



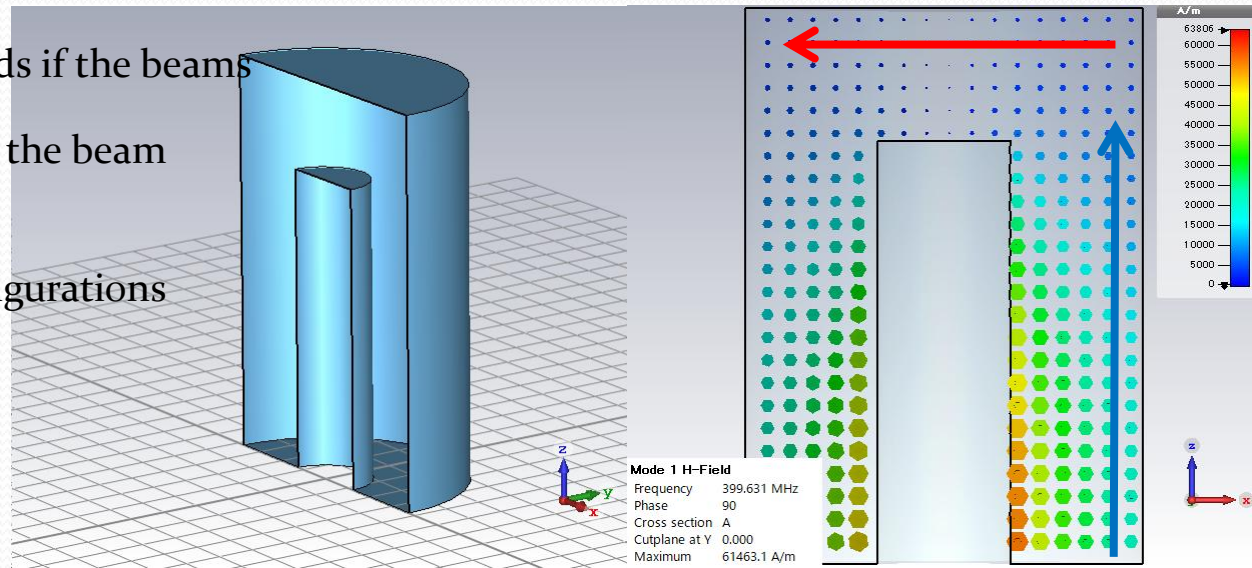
Compact crab cavities for LHC

- Required for LHC application
- Required voltage: 3MV
- Frequency: 400 MHz
- Operating temperature: 2K
- Narrow space between beam pipes
 - Nominal beam separation: 194 mm
 - Beam pipe diameter: 84 mm
- Can not use pill box type cavities
 - 800MHz cavity clears size problem but the high frequency gives non-linear kick



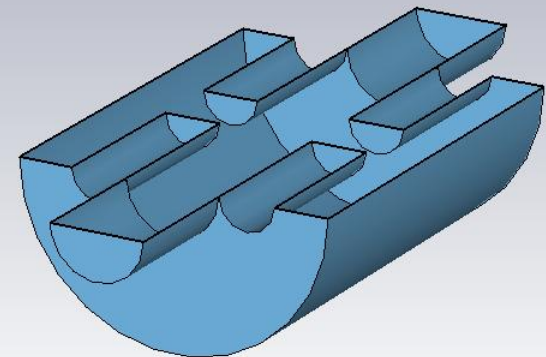
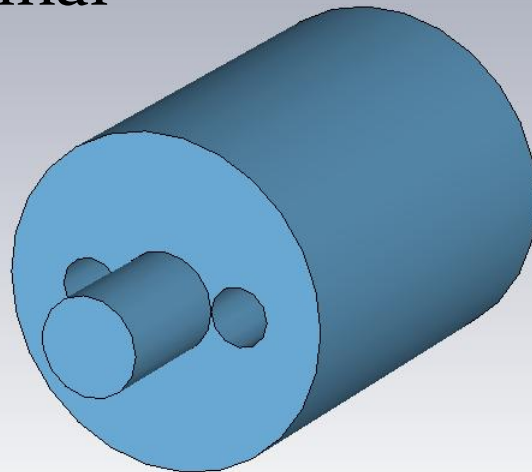
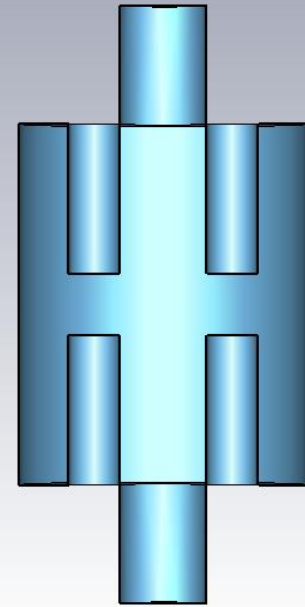
Quarter wave resonator

- Compact cavity
- Used for low beta beam accelerations
 - Low frequency application
 - If you make a 200MHz pill box cavity, its diameter should be 1 m.
- Maximum E field exists at the open end of an inner conductor
 - Deflection by the E-field if the beams pass along the red line
- Maximum H field exists at the short end of an inner conductor
 - Deflection by the H fields if the beams pass along the blue line
 - E fields also contributes the beam deflection
- Beam acceleration
 - 2 or 4 symmetrical configurations avoid acceleration



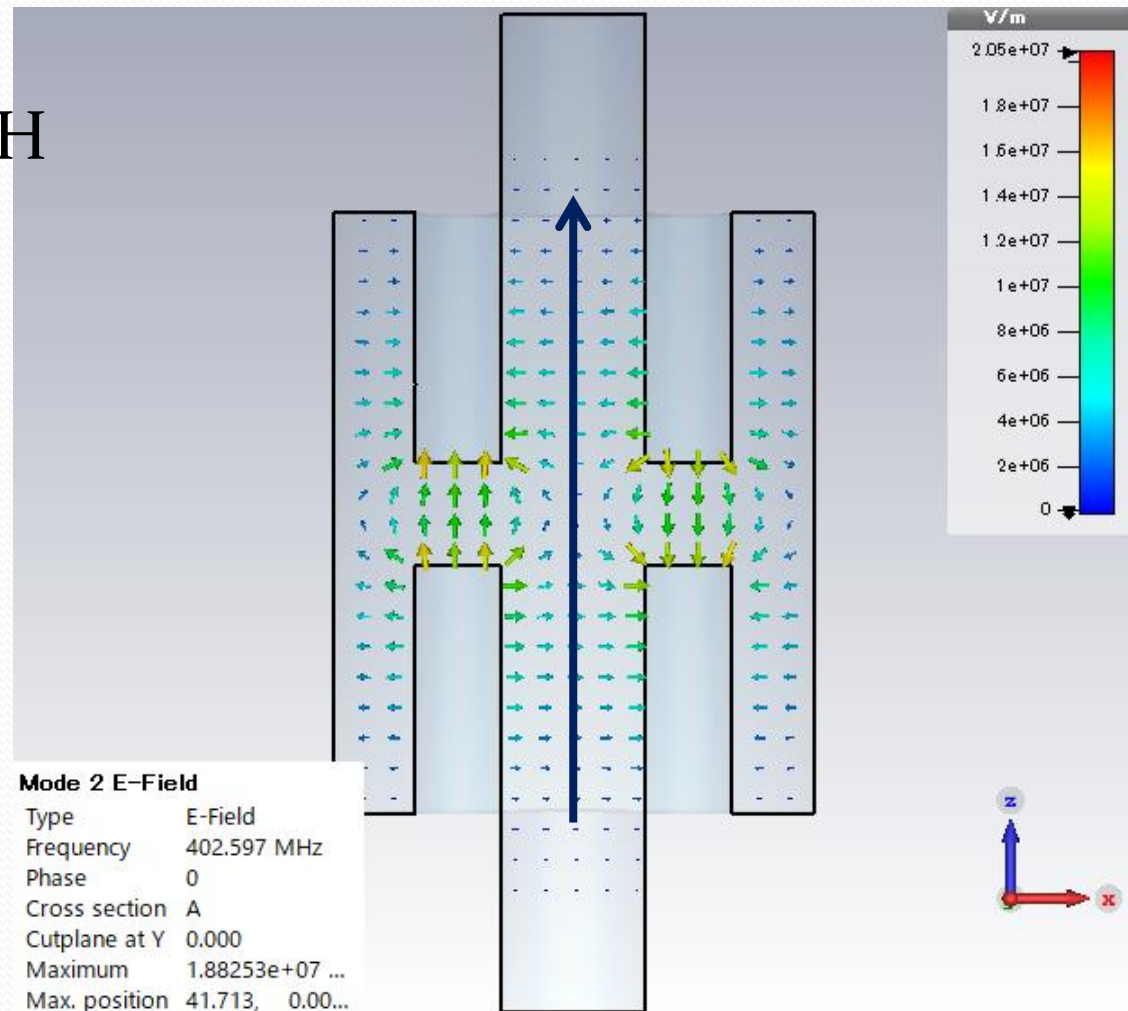
Four QWR configuration 4RCC

- Consists of four QWRs
- Use the E and B fields for beam deflection
- $R/Q \sim 1000$
- No longitudinal acceleration
- Not FM



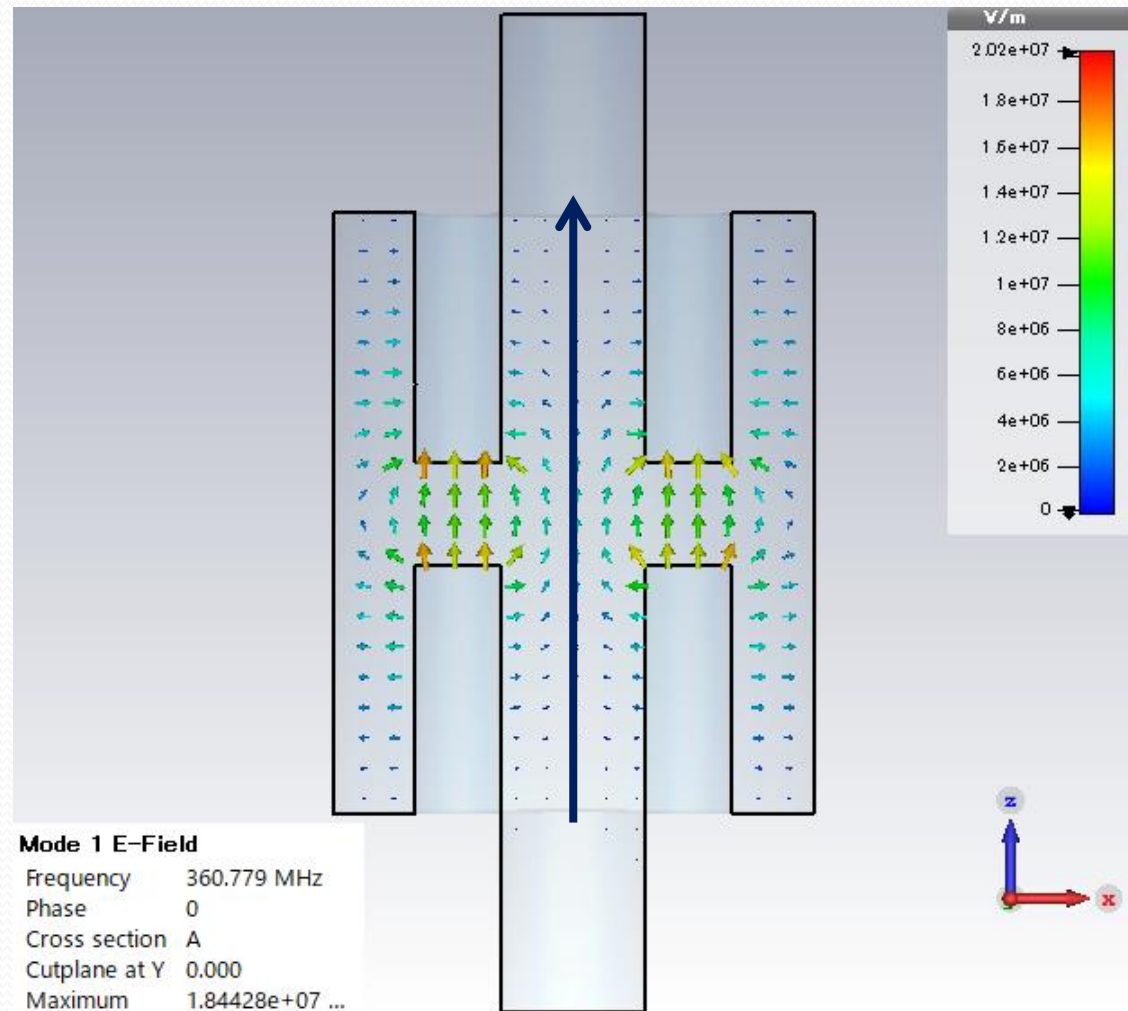
Crabbing mode

- 400 MHz
- Deflecting by E and H fields
- No longitudinal acceleration
- $R/Q \sim 1000$



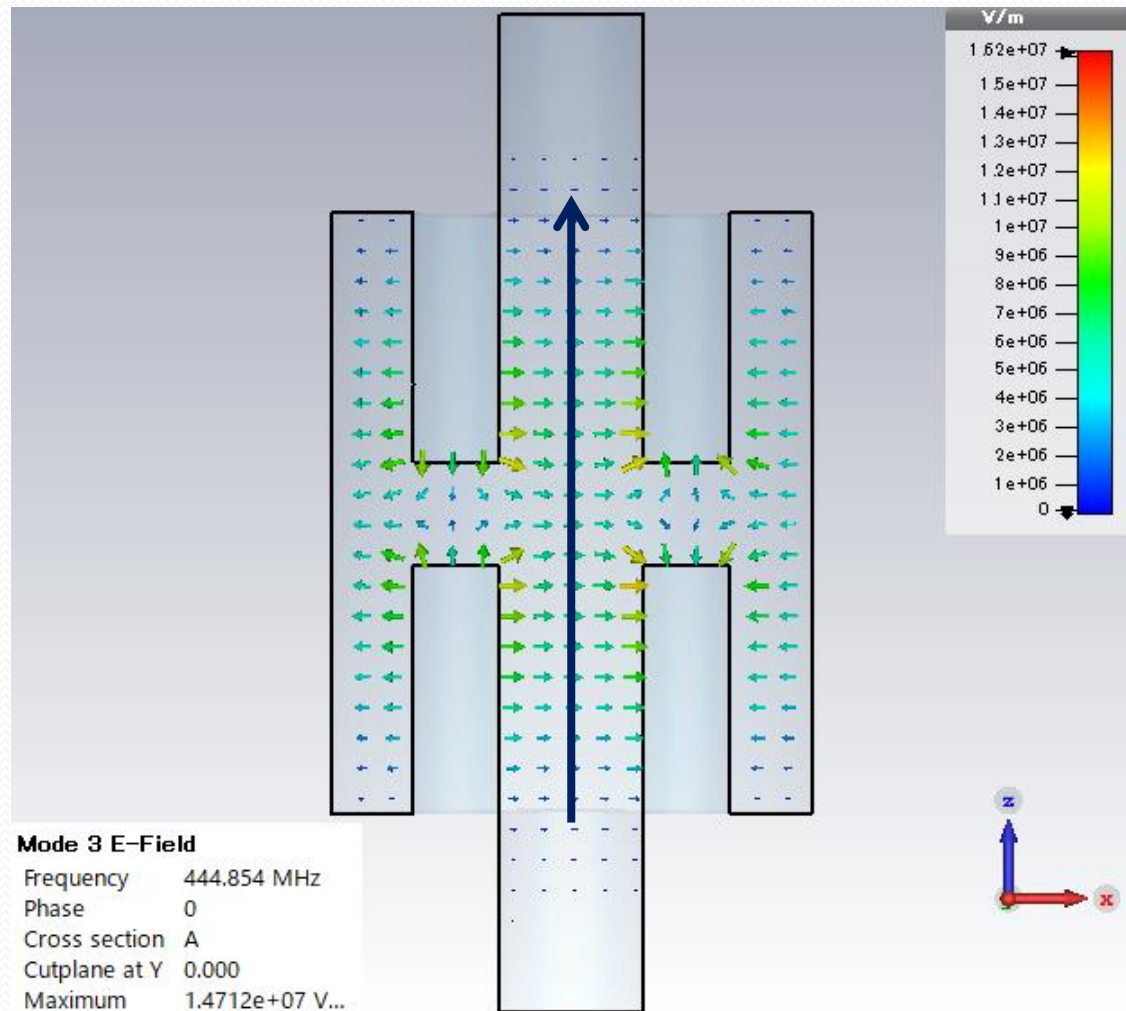
Fundamental mode

- 360 MHz
- Accelerating mode



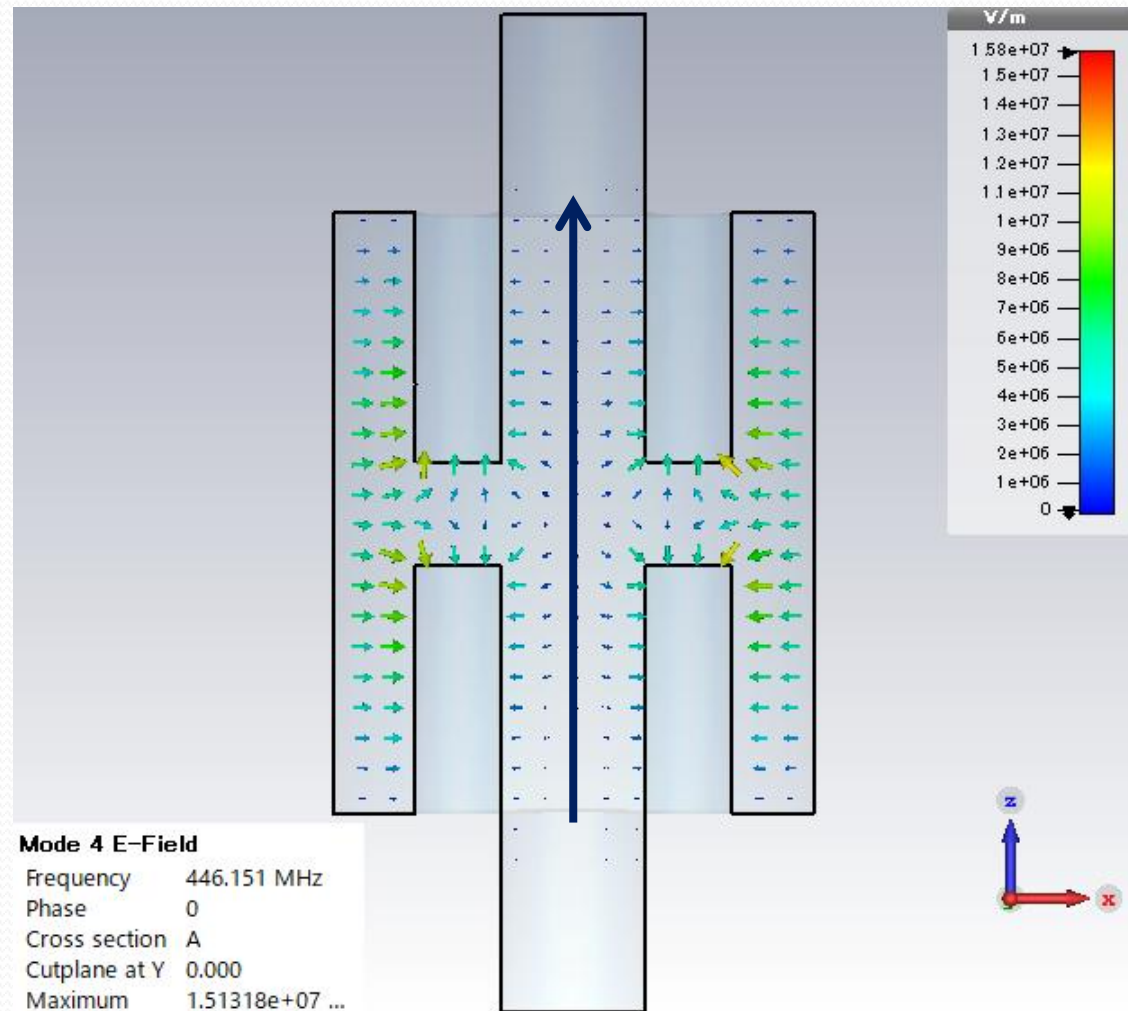
1st HOM

- 445 MHz
- Deflecting mode



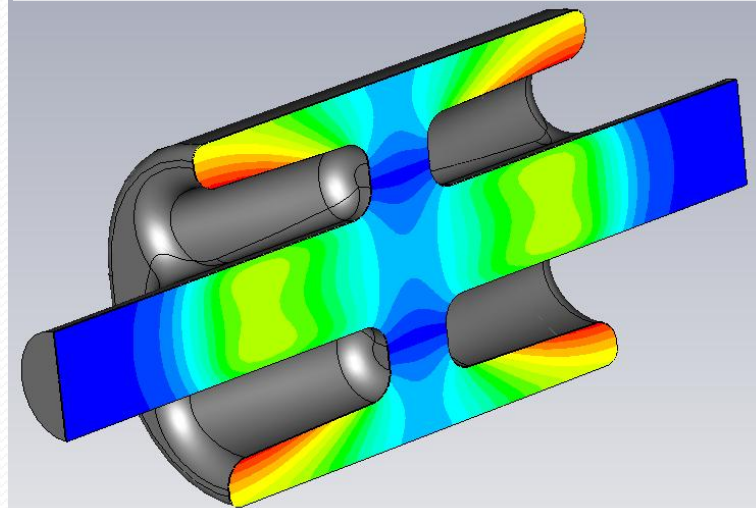
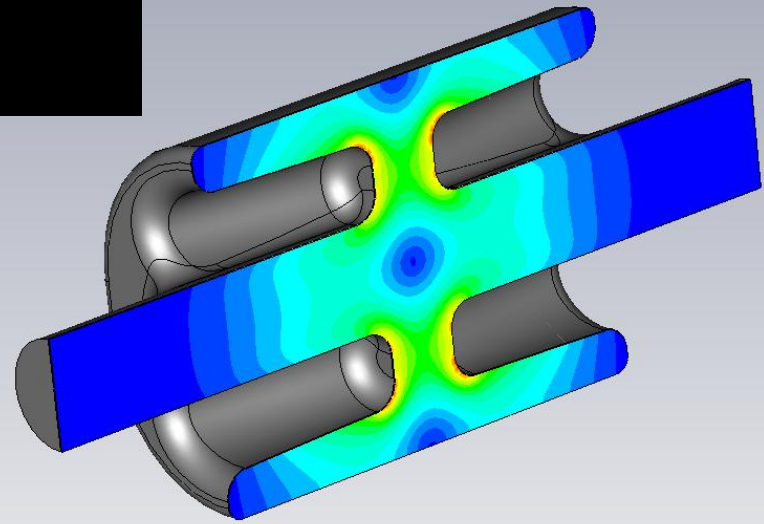
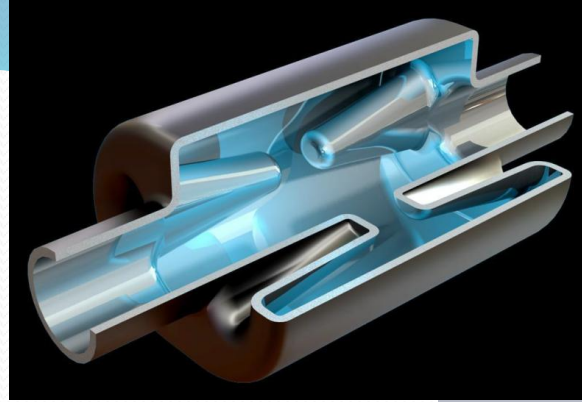
2nd HOM

- 446 MHz
- Accelerating mode



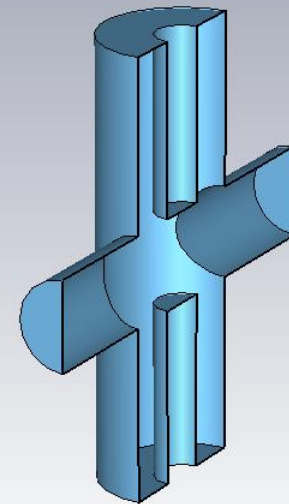
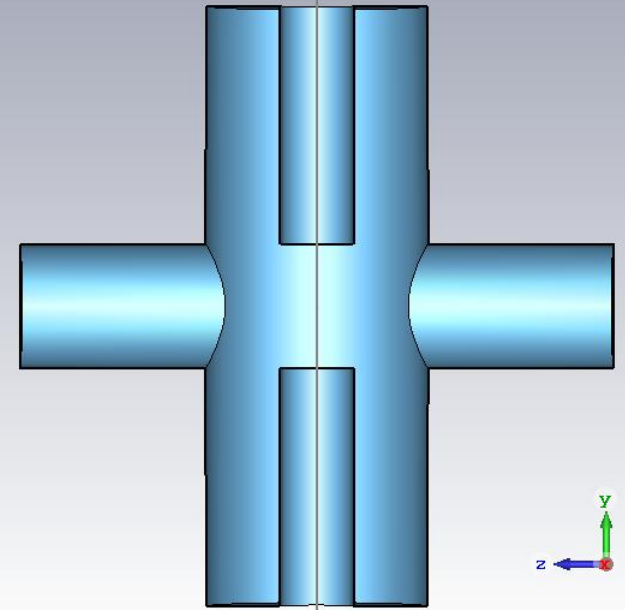
4RCC(UK)

- 4-rod crab cavity
- ULANC
- Cavity size: $143 \times 118 \text{ mm}^2$
- Cavity length: 500 mm
- Peak E field: 32 MV/m
- Peak B-Field 61 mT
- R/Q: 915
- FM: 370 MHz
- Not fundamental mode



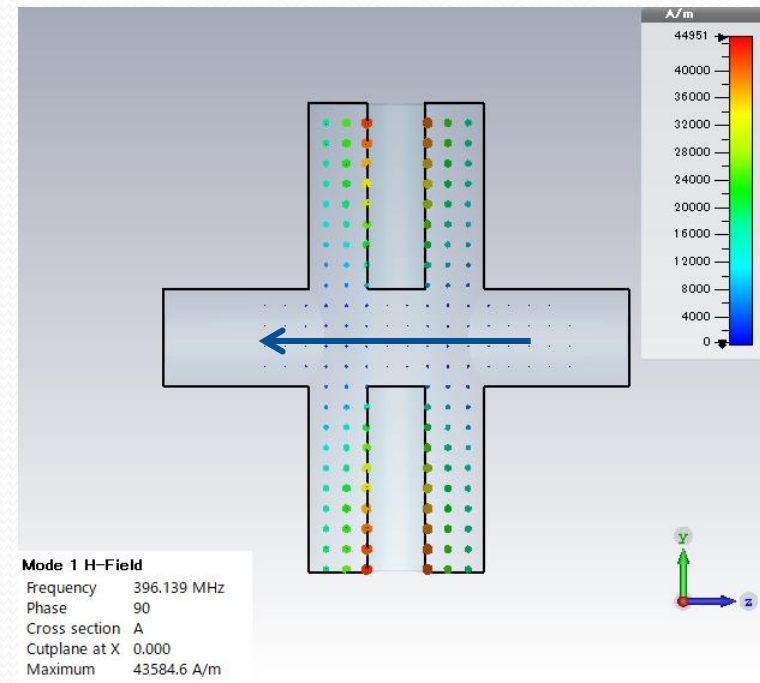
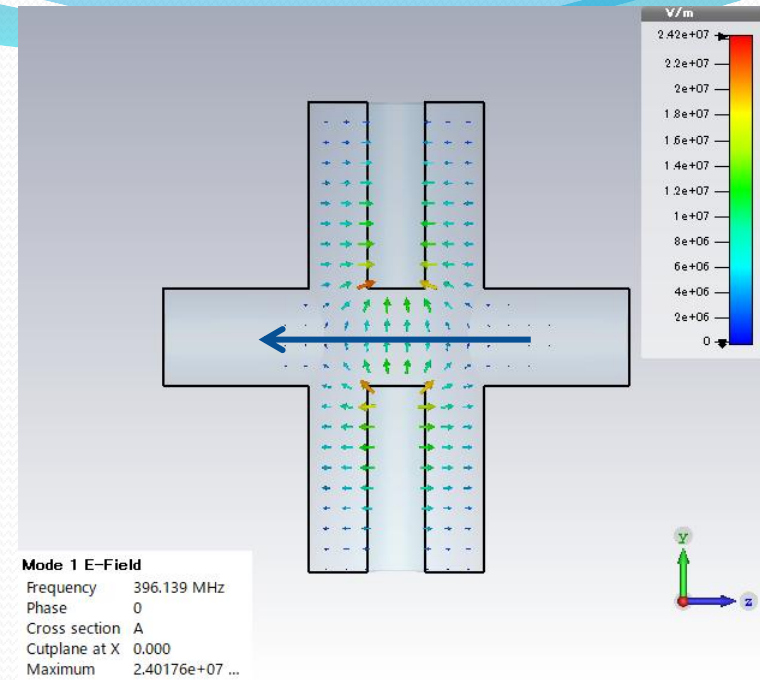
Two QWR configuration DQW

- Consists two QWRs
- Mainly use the E-field for beam kick
- No longitudinal acceleration
- FM



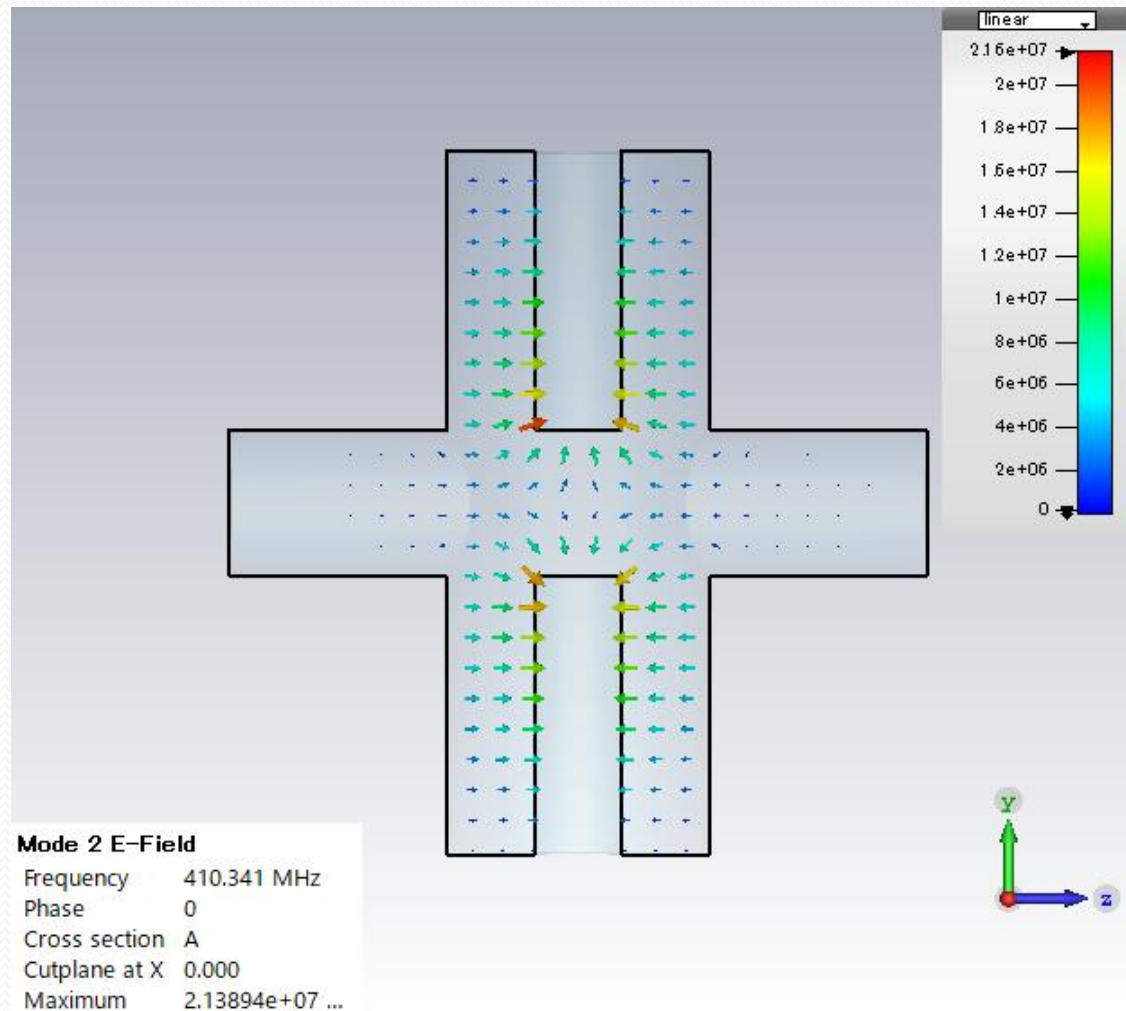
Crabbing mode

- 400 MHz
- Deflecting mainly by the E field
- No longitudinal acceleration



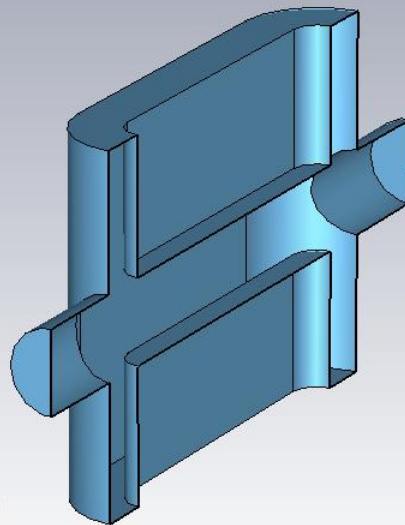
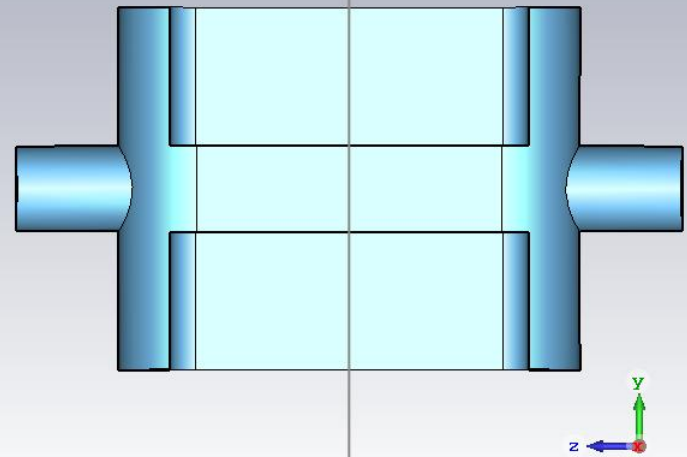
HOM

- 410 MHz
 - Near the crabbing frequency
 - Well separate after longitudinal length optimization
- Accelerating mode

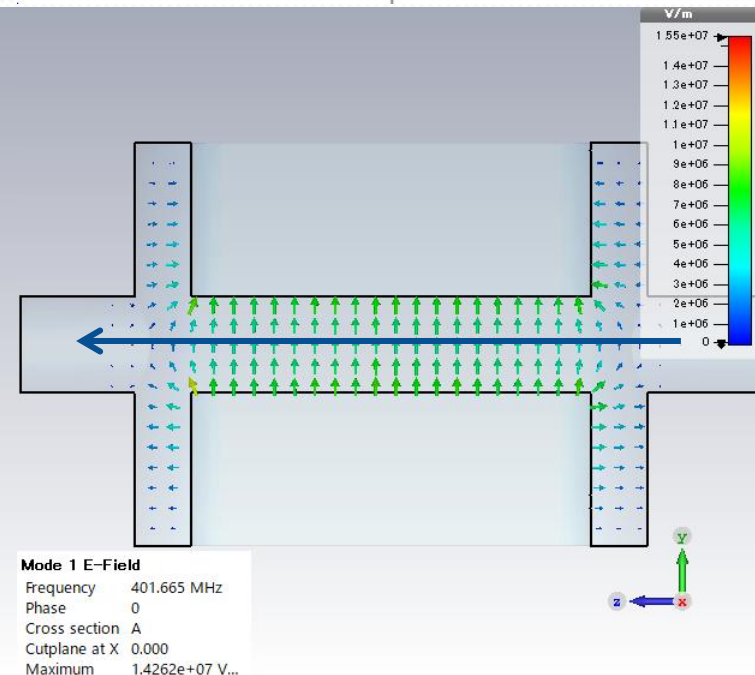


Longitudinal length

- To obtain larger kick voltage
 - Deflecting area has to be increased
 - Optimize longitudinal length $\sim \lambda/2$
- $R/Q \sim 500$
- HOM: 457 MHz

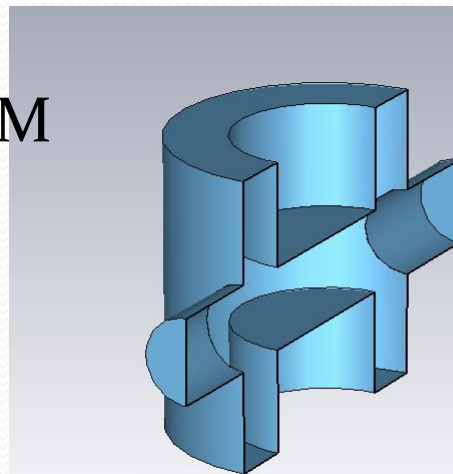


component1.bp
Material Vacuum
Type Normal
Epsilon 1
Mu 1

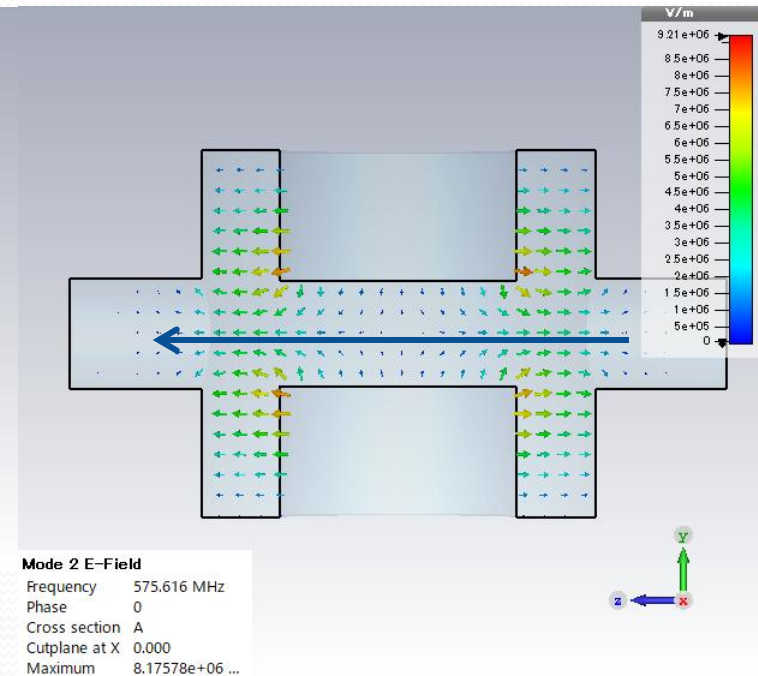
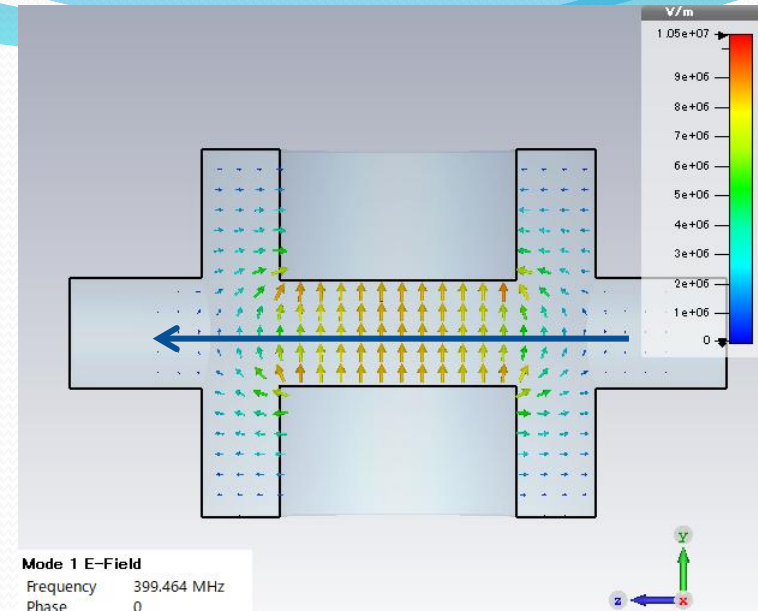


Transverse length

- Increase of transverse length
 - Results in shorter inner conductor
 - 160 to 100 mm
- Increasing HOM frequency
 - 575 MHz
 - Well above FM

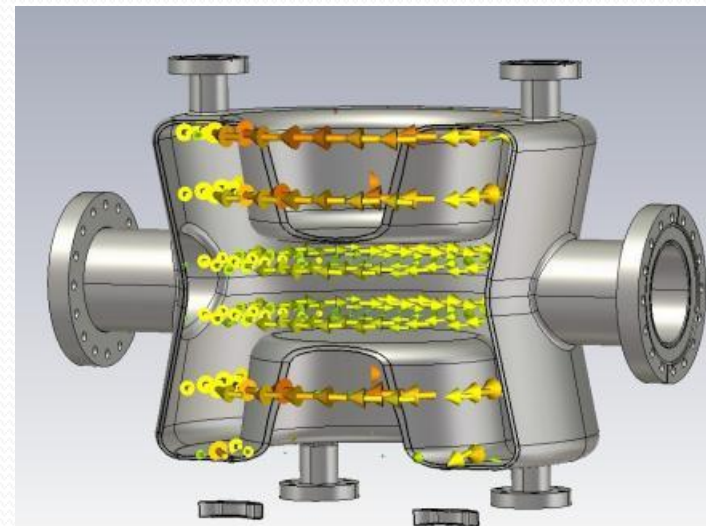
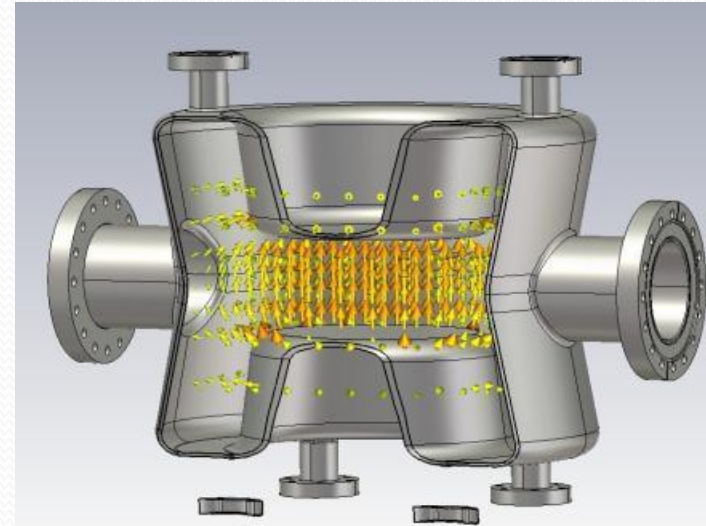
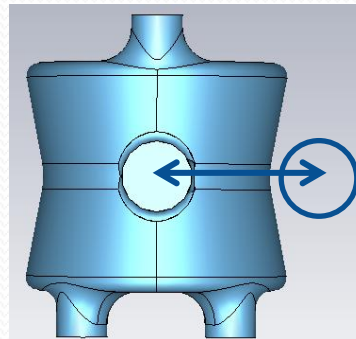


component1:outer
Material Vacuum
Type Normal
Epsilon 1
Mu 1



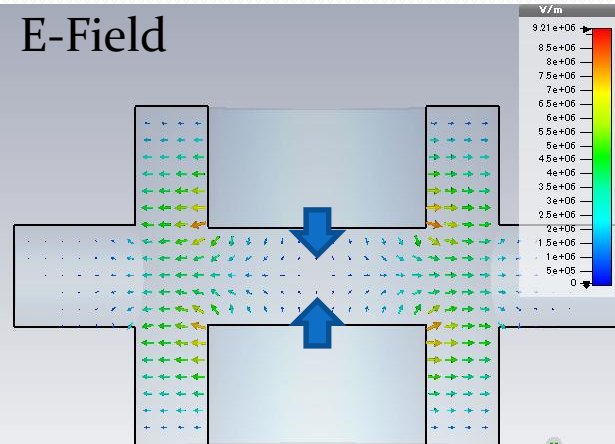
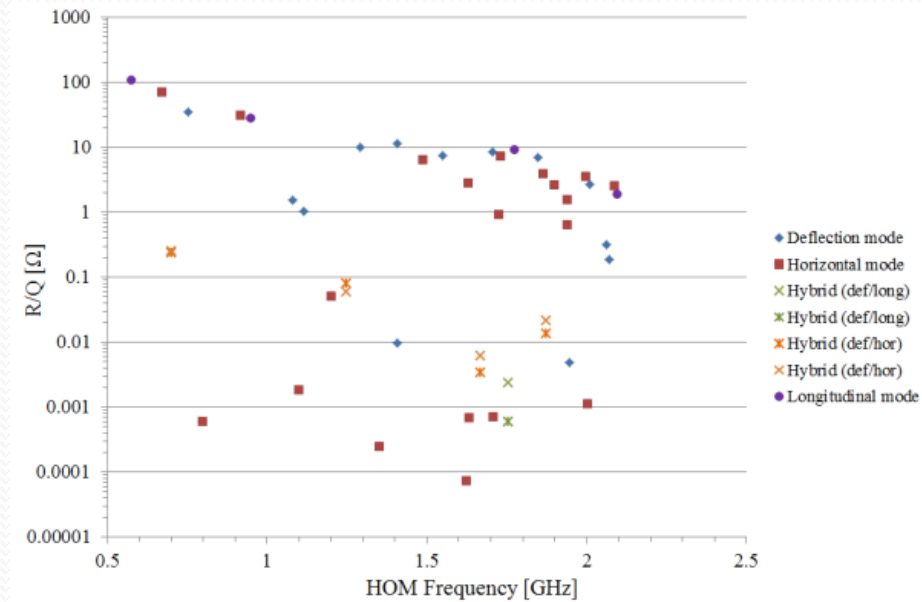
DQW

- Double quarter wave cavity
- BNL
- Vertical kick
- FM
- Cavity size: $142 \times 122 \text{ mm}^2$
 - Need slim waist to provide space for the other BP
- Cavity length: 384 mm
- Peak E field: 43 MV/m
- Peak B-Field 61 mT
- R/Q: 345
- HOM: 575 MHz

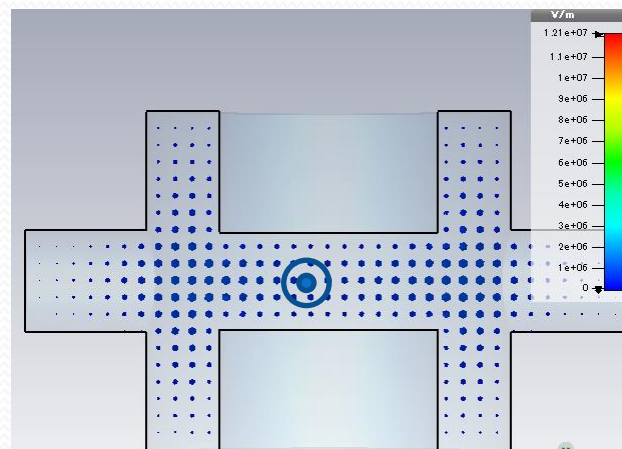


HOMs

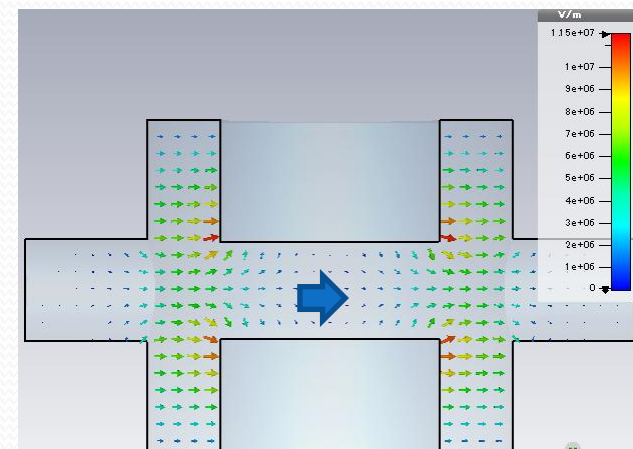
- HOMs are well separate from FM
- Above 550 MHz



Mode 2 E-Field
 Frequency 575.616 MHz
 Phase 0
 Cross section A
 Cutplane at X 0.000
 Maximum 8.17578e+06 ...



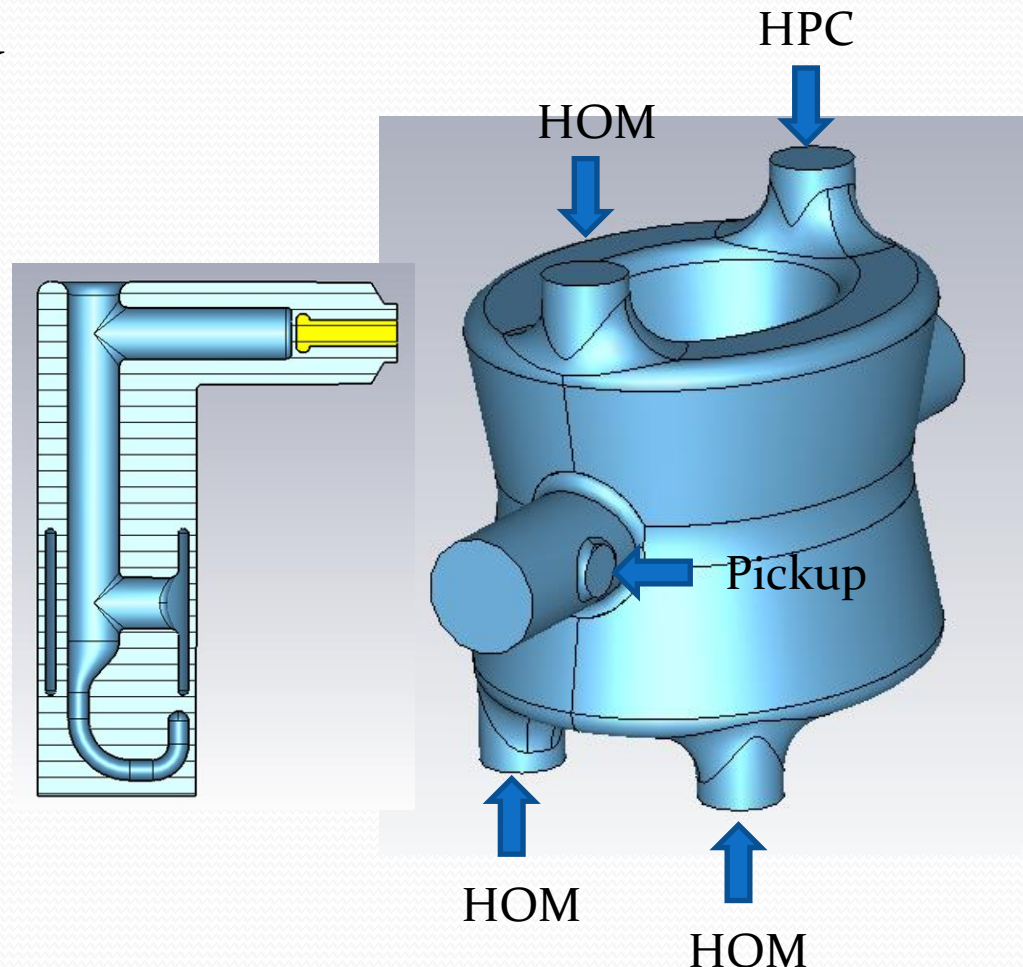
Mode 3 E-Field
 Frequency 702.144 MHz
 Phase 0
 Cross section A
 Cutplane at X 0.000
 Maximum 1.20825e+06 ...



Mode 4 E-Field
 Frequency 710.519 MHz
 Phase 0
 Cross section A
 Cutplane at X 0.000
 Maximum 1.12688e+07 ...

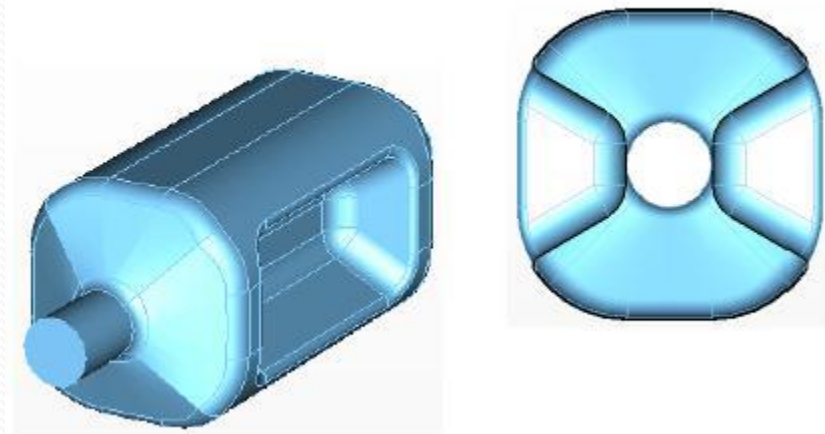
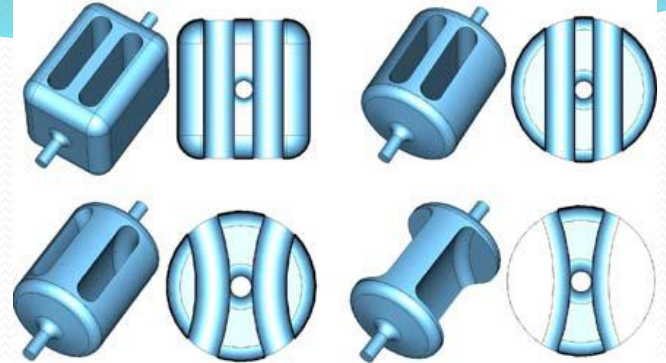
HOM damping

- HOMs are extracted by couplers
 - Hook-type coupler
 - Band stop filter at 400MHz
 - L-type high pass filter above 570MHz

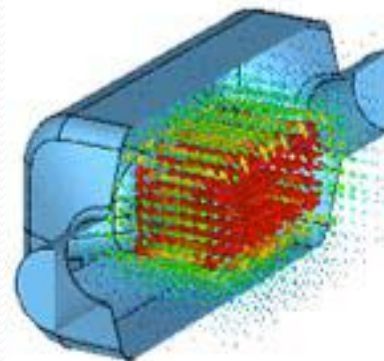


RFD

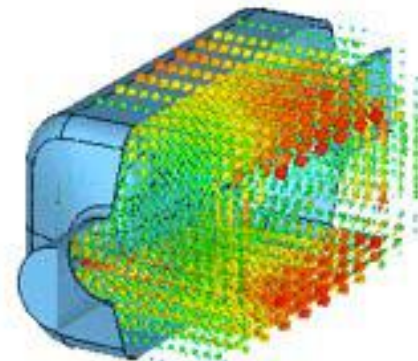
- RF deflector
 - Same topology
 - Evolved from the parallel bar geometry
- ODU/JLAB/SLAC
- Horizontal kick
- FM
- Cavity size: $147 \times 147 \text{ mm}^2$
- Cavity length: 600 mm
- Peak E field: 33 MV/m
- Peak B-Field 56 mT
- R/Q: 287
- HOM: 580 MHz



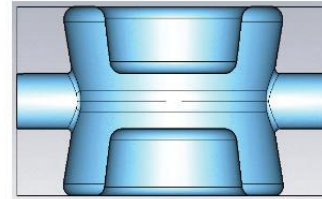
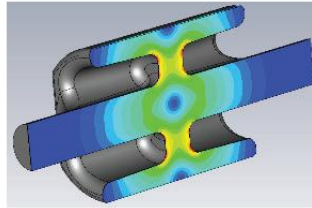
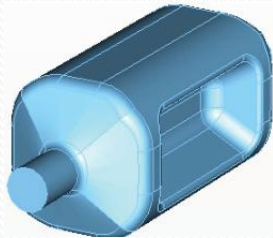
E Field



H Field



Compact CC



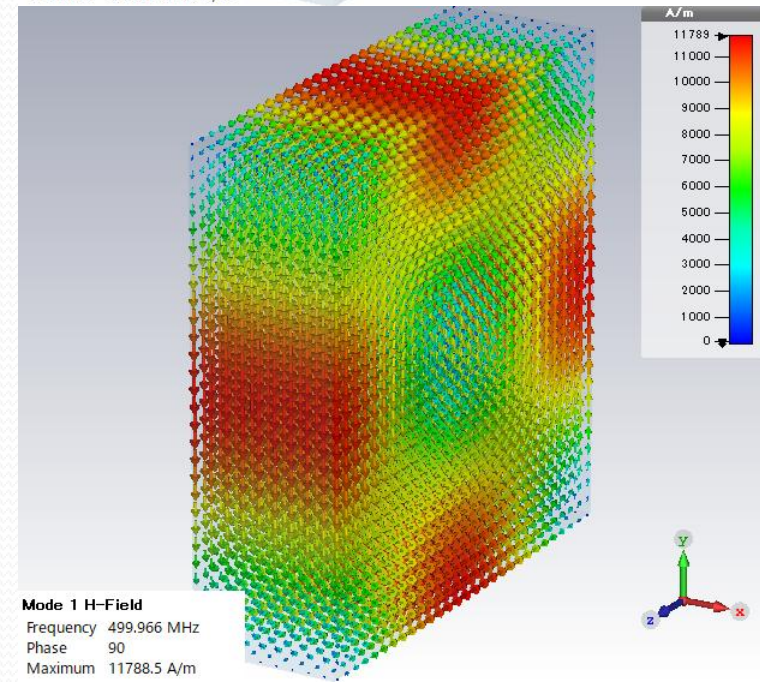
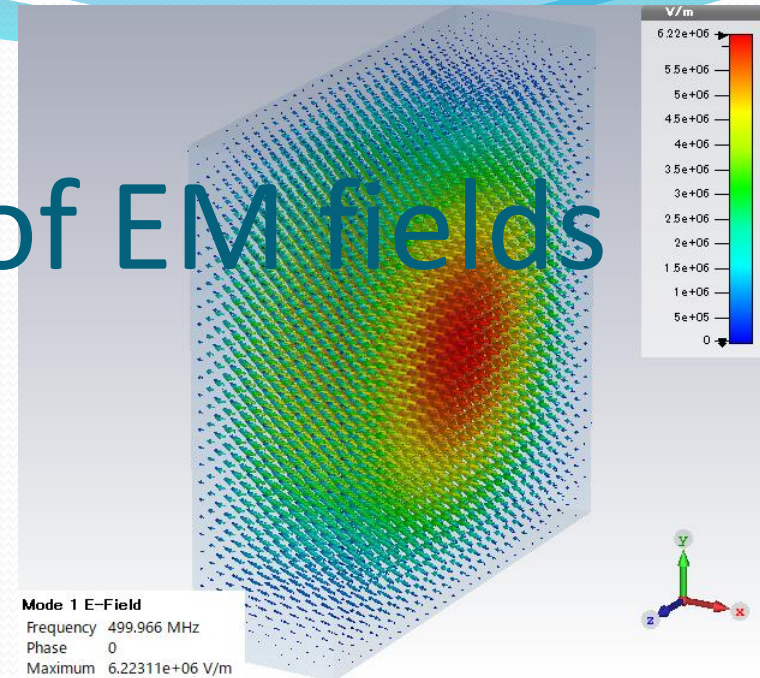
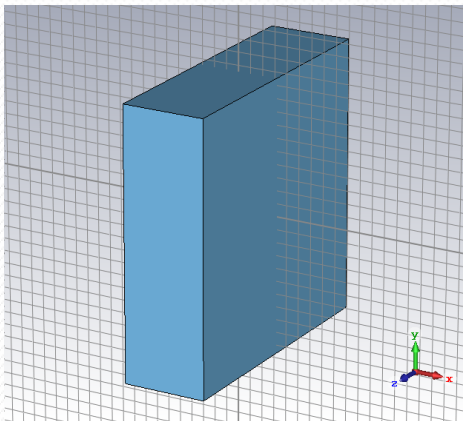
400MHz, 3MV	RFD	4RCC	DQW
Cavity radius (mm)	147.5	143/118	142/122
Cavity length (mm)	597	500	380
Beam pipe radius (mm)	42	42	42
Peak E-field (MV/m)	33	32	47
Peak B-field (mT)	56	61	71
R/Q (Ω)	287	915	318
Nearest mode (MHz)	584	370	575

RFD and DQW were chosen for LHC CC



Analytical solution of EM fields

- Consider EM fields in a rectangular cavity
- Obtain analytical solutions of the TE₀₁₁ mode
- Calculate its kick voltage



Electro-magnetic fields

$$E_x = iE_0 \frac{kk_y}{k_c^2} \cos k_x x \sin k_y y \sin k_z z$$

$$E_y = -iE_0 \frac{kk_x}{k_c^2} \sin k_x x \cos k_y y \sin k_z z$$

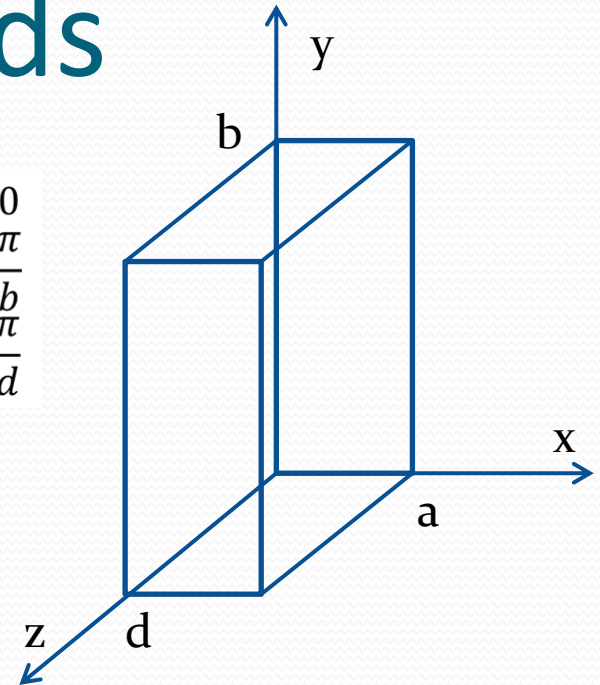
$$E_z = 0$$

$$H_x = -\frac{E_0}{c\mu} \frac{k_x k_z}{k_c^2} \sin k_x x \cos k_y y \cos k_z z$$

$$H_y = -\frac{E_0}{c\mu} \frac{k_y k_z}{k_c^2} \cos k_x x \sin k_y y \cos k_z z$$

$$H_z = \frac{E_0}{c\mu} \cos k_x x \cos k_y y \sin k_z z$$

$$\begin{aligned} k_x &= 0 \\ k_y &= \frac{\pi}{b} \\ k_z &= \frac{\pi}{d} \end{aligned}$$



$$CB_x = -E_0 \frac{k_x k_z}{k_c^2} \sin k_x x \cos k_y y \cos k_z z$$

$$CB_y = -E_0 \frac{k_y k_z}{k_c^2} \cos k_x x \sin k_y y \cos k_z z$$

$$CB_z = E_0 \cos k_x x \cos k_y y \sin k_z z$$

Integration of analytical solution

- Calculate kick voltage for the TE mode in a rectangular cavity
 - Electromagnetic fields are given
 - A particle passes at $x=a/2, y=b/2$
 - Time factor $e^{i\omega t}$

$$E_x = iE_0 \frac{kk_y}{k_c^2} \sin k_y y \sin k_z z$$
$$CB_y = -E_0 \frac{k_y k_z}{k_c^2} \sin k_y y \cos k_z z$$

$$E_x = iE_0 \frac{kk_y}{k_c^2} \sin k_z z$$
$$CB_y = -E_0 \frac{k_y k_z}{k_c^2} \cos k_z z$$

$$E_x = iE_0 \frac{kk_y}{k_c^2} \sin k_z z \quad e^{i\omega t}$$
$$CB_y = -E_0 \frac{k_y k_z}{k_c^2} \cos k_z z \quad e^{i\omega t}$$

Integration of analytical solution

- Calculate kick voltage for the TE mode in a rectangular cavity
 - Real part
 - Imaginary part

$$E_x = -E_0 \frac{k k_y}{k_c^2} \sin k_z z \sin kz$$
$$CB_y = -E_0 \frac{k_y k_z}{k_c^2} \cos k_z z \cos kz$$

$$E_x = E_0 \frac{k k_y}{k_c^2} \sin k_z z \cos kz$$
$$CB_y = -E_0 \frac{k_y k_z}{k_c^2} \cos k_z z \sin kz$$

Integration of analytical solution

- Calculate kick voltage for the TE mode in a rectangular cavity
 - Kick voltage
 - Field integration along beam passage
 - Real part=0
 - Imaginary part=0
 - Kick voltage=0

$$|V_x| = 0$$

$$V_x = \int_0^d (E_x - cB_y) e^{ikz} dz$$

$$\int_0^d \sin k_z z \cos kz dz = -\frac{k_z}{k_c^2} (1 + \cos kd)$$

$$\int_0^d \cos k_z z \sin kz dz = \frac{k}{k_c^2} (1 + \cos kd)$$

$$\int_0^d \sin k_z z \sin kz dz = \frac{k_z}{k_c^2} \sin kd$$

$$\int_0^d \cos k_z z \cos kz dz = \frac{k}{k_c^2} \sin kd$$

$$\text{Re}V_x = \text{Re} \int_0^d (E_x - cB_y) e^{ikz} dz = 0$$

$$\text{Im}V_x = \text{Im} \int_0^d (E_x - cB_y) e^{ikz} dz = 0$$